

Negative Komar Mass in regular stationary spacetimes

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Plan of the Talk

- 1 Introduction**
- 2 The Komar mass**
- 3 Journey into the realm of negative Komar mass**
- 4 Summary**

1. Introduction

- We consider a self-gravitating system consisting of a uniformly rotating, homogeneous perfect fluid ring and a central object, being either a black hole or a disk of dust.

- Axisymmetry and stationarity are described by Killing vectors

$$\eta^i \quad \text{and} \quad \xi^i$$

- Line element in Weyl-Lewis-Papapetrou coordinates:

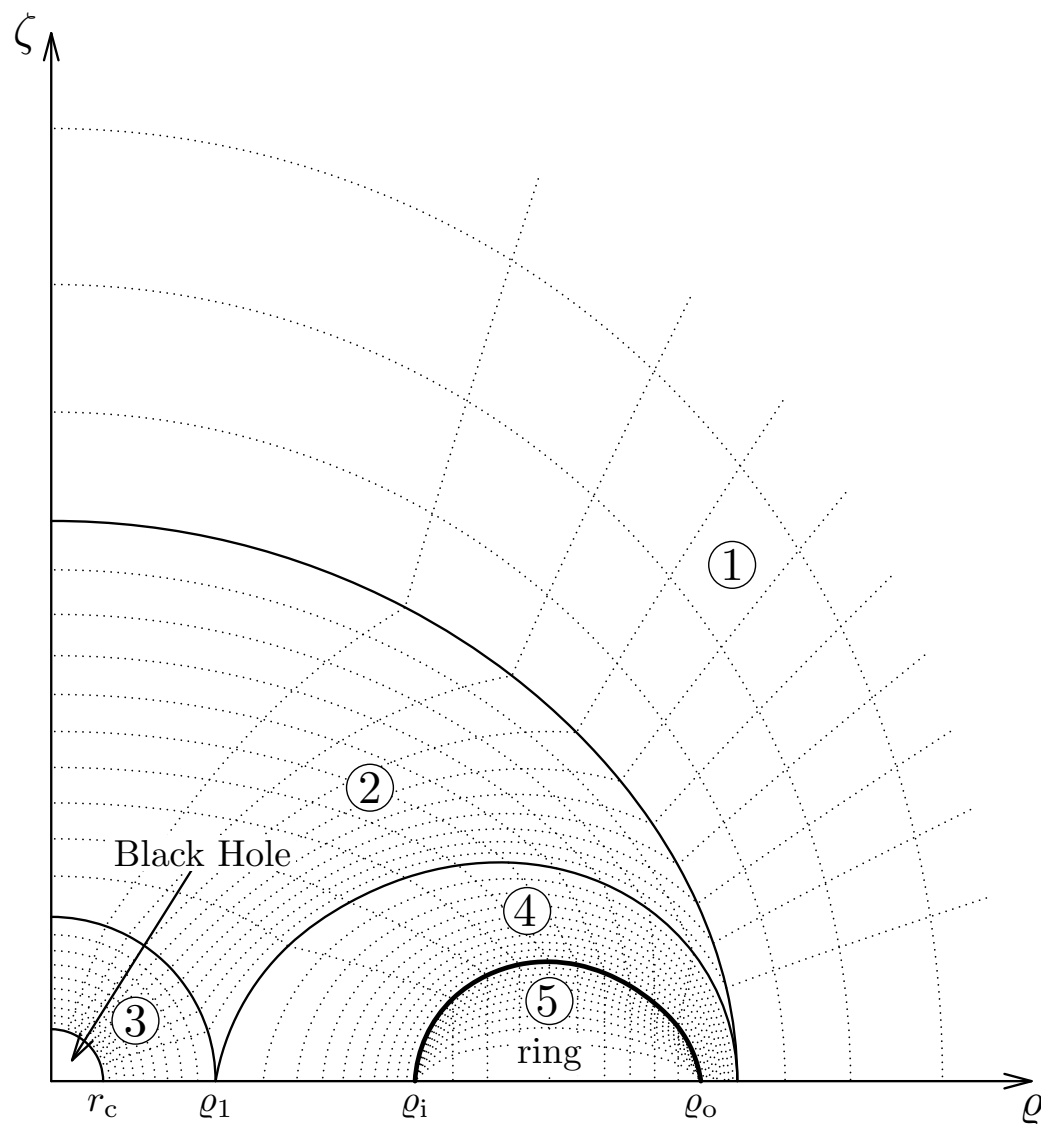
$$ds^2 = -e^{2\nu} dt^2 + \varrho^2 B^2 e^{-2\nu} (d\varphi - \omega dt)^2 + e^{2\lambda} (d\varrho^2 + d\zeta^2)$$

- For the metric functions ν, B, ω, λ we solve the corresponding free boundary value problem with spectral methods.

- The presence of the ring can affect the properties of the central object drastically.

- We illustrate the ring's influence by tracing paths along which the 'Komar' mass of the central object becomes negative.

1. Introduction



The division of the ρ - ζ plane into the domains used in the spectral methods ($\rho_i/\rho_o = 0.56$ and $r_c/\rho_o = 0.08$).

2. The Komar mass (1)

- Poisson equation in Newtonian gravity

$$\nabla \cdot (\nabla U) = 4\pi\mu$$

- A mass can be assigned to any subregion $V \subseteq \mathbb{R}^3$

$$M(V) = \int_V \mu d^3x = \frac{1}{4\pi} \oint_{\partial V} \nabla U d\vec{f}$$

- Consequences:

1) $M(V) = 0$ if V is a vacuum region with $\mu = 0$

2) $M_{\text{total}} = M(\mathbb{R}^3) = - \lim_{r \rightarrow \infty} (rU)$

2. The Komar mass (2)

- Specific Einstein equation in axisymmetry and stationarity:

$$\nabla \cdot \left(B \nabla \nu - \frac{\omega}{2} \varrho^2 B^3 e^{-4\nu} \nabla \omega \right) = 4\pi \tilde{\mu}(\mu, p; \lambda, \nu, B, \omega)$$

- A 'Komar' mass can be assigned to any subregion $V \subseteq \mathbb{R}^3$

$$M(V) = \int_V \tilde{\mu} d^3x = \frac{1}{4\pi} \oint_{\partial V} \left(B \nabla \nu - \frac{\omega}{2} \varrho^2 B^3 e^{-4\nu} \nabla \omega \right) d\vec{f}$$

- Consequences:

$$1) \quad M(V) = 0 \quad \text{if } V \text{ is a vacuum region with } \mu = 0$$

$$2) \quad M_{\text{ADM}} = M(\mathbb{R}^3) = - \lim_{r \rightarrow \infty} (r\nu)$$

- Define the Komar mass of a black hole as surface integral over an arbitrary boundary ∂V where V contains only the black hole.

2. The Komar mass (3)

- **Question:** Is the black hole's Komar mass always positive ?
- Analysis by means of the 'Smarr' formula (Bardeen, Carter):

$$M_h = \frac{\kappa}{4\pi} A_h + 2\Omega_h J_h$$

- The surface gravity κ and horizon area A_h are always positive, but the product can approach zero.
- The ring can cause a 'frame dragging' of the black hole such that its angular velocity Ω_h and angular momentum J_h assume different signs.
- Requirement: highly relativistic rotating rings, characterized by a large ergosphere (a portion of space in which $\xi^i \xi_i > 0$).
- Can the negative term $2\Omega_h J_h$ dominate over $\kappa A_h / (4\pi)$?
- Answer: **Yes!** The Komar mass of such black holes is **negative**.

2. The Komar mass (4)

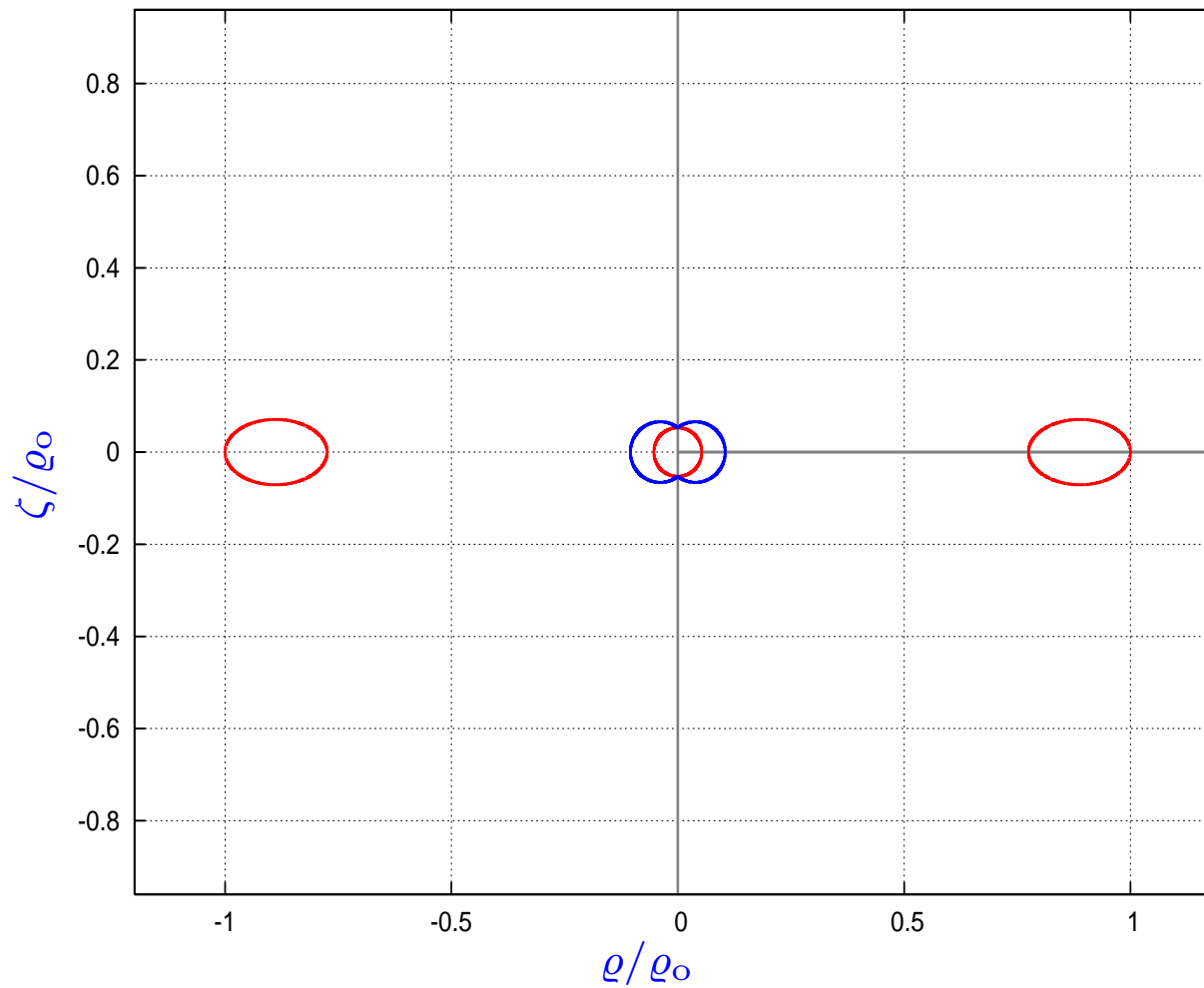
- **Question:** Is the central disk's Komar mass always positive ?
- Analysis by means of the Disk-'Smarr' formula (Bardeen):

$$M_d = e^{V_0} M_0 + 2\Omega_d J_d$$

- The redshift $Z_d = e^{-V_0} - 1$ and the *baryonic* mass M_0 are always positive, but the product $e^{V_0} M_0$ can approach zero.
- Again, a 'frame dragging' caused by the ring can lead to different signs of the disk's angular velocity Ω_d and its angular momentum J_d .
- Requirement: highly relativistic rotating rings, characterized by a large ergosphere (a portion of space in which $\xi^i \xi_i > 0$).
- Can the negative term $2\Omega_d J_d$ dominate over $e^{V_0} M_0$?
- **Answer: Yes!** The Komar mass of such dust disks is **negative**.

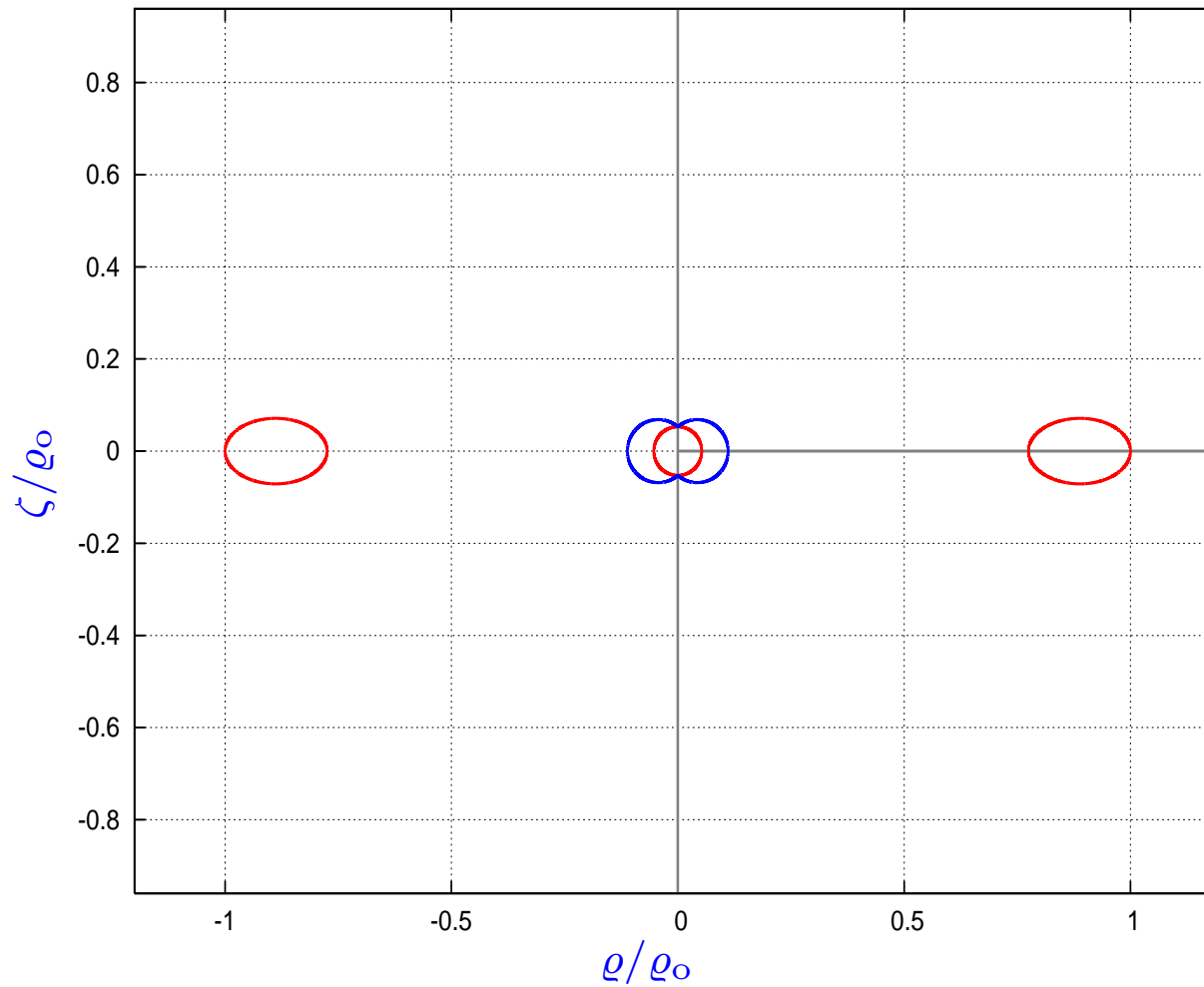
3. Journey

$$M_h/M_r = 0.89, Z_r = 0.65$$



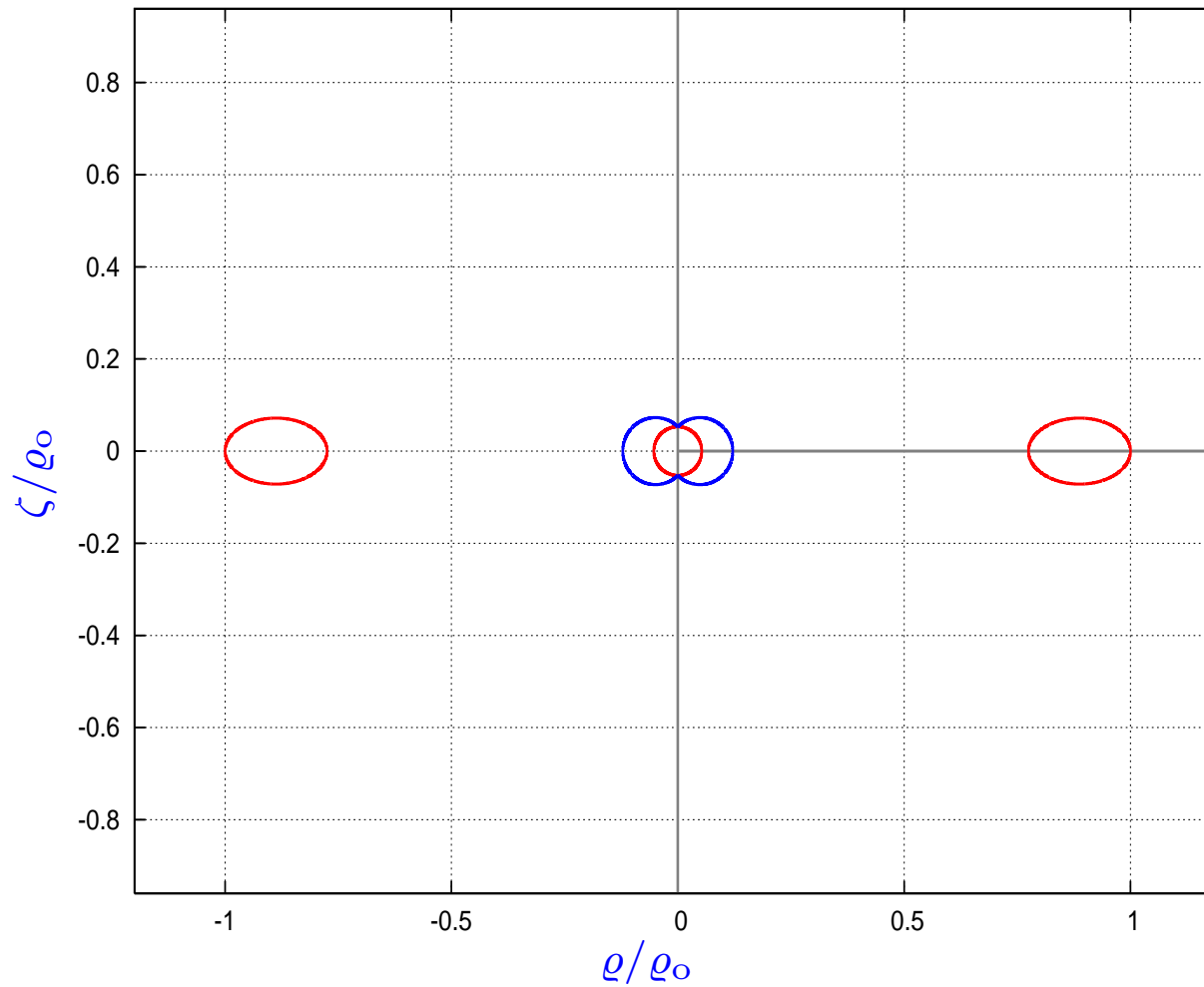
3. Journey

$$M_h/M_r = 0.78, Z_r = 0.75$$



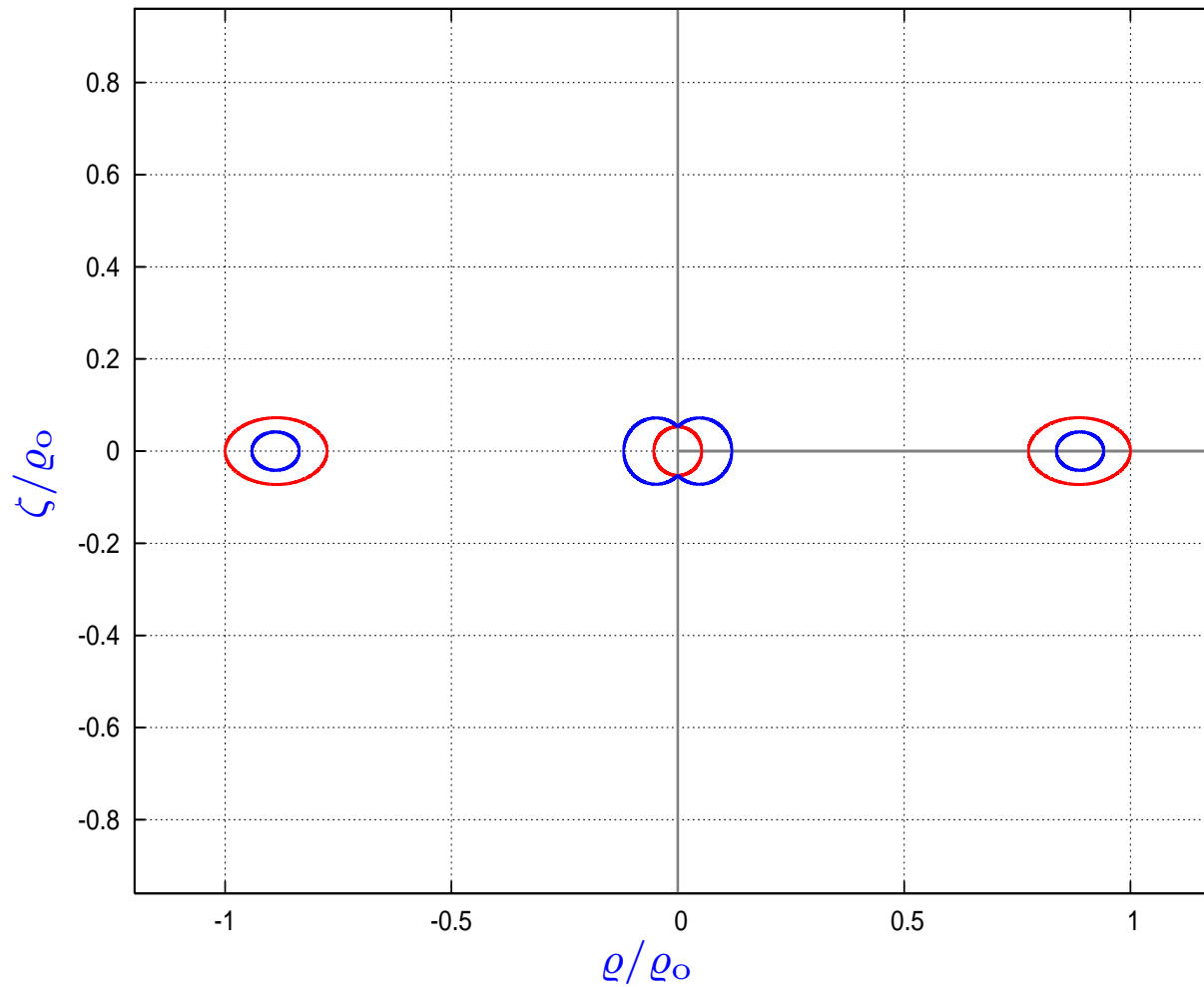
3. Journey

$$M_h/M_r = 0.53, Z_r = 1.1$$



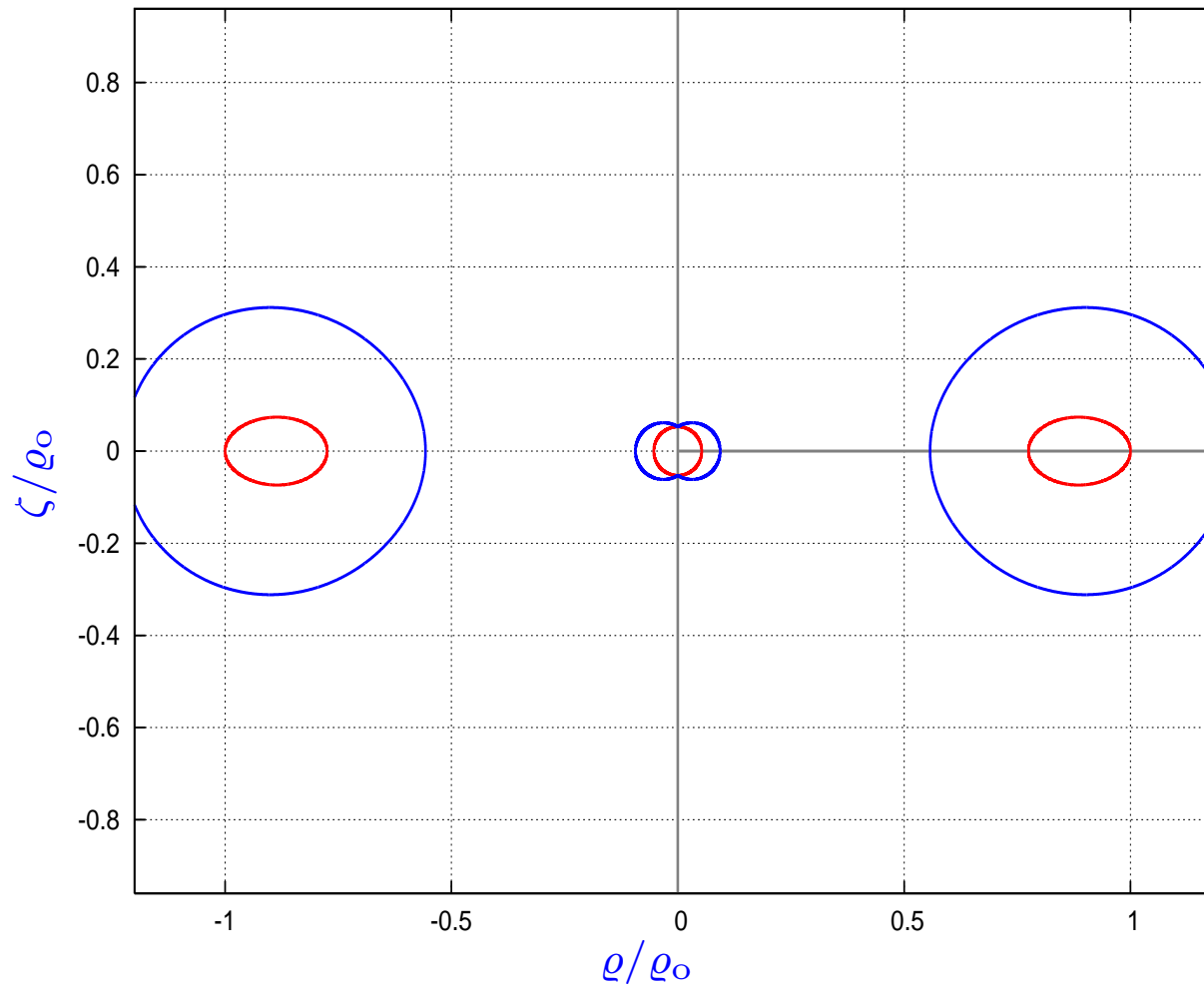
3. Journey

$$M_h/M_r = 0.33, Z_r = 1.6$$



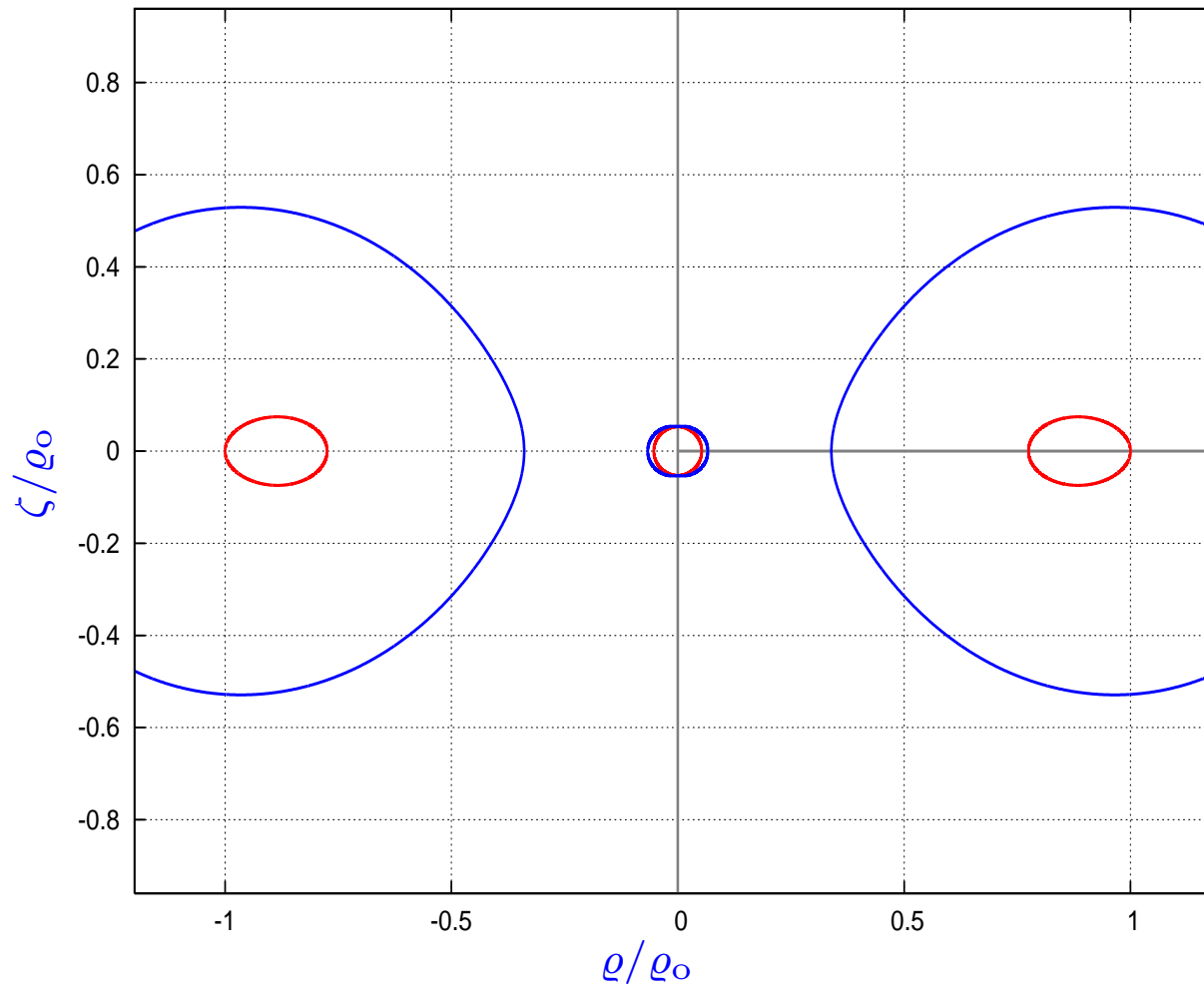
3. Journey

$$M_h/M_r = 0.16, Z_r = 2.7$$



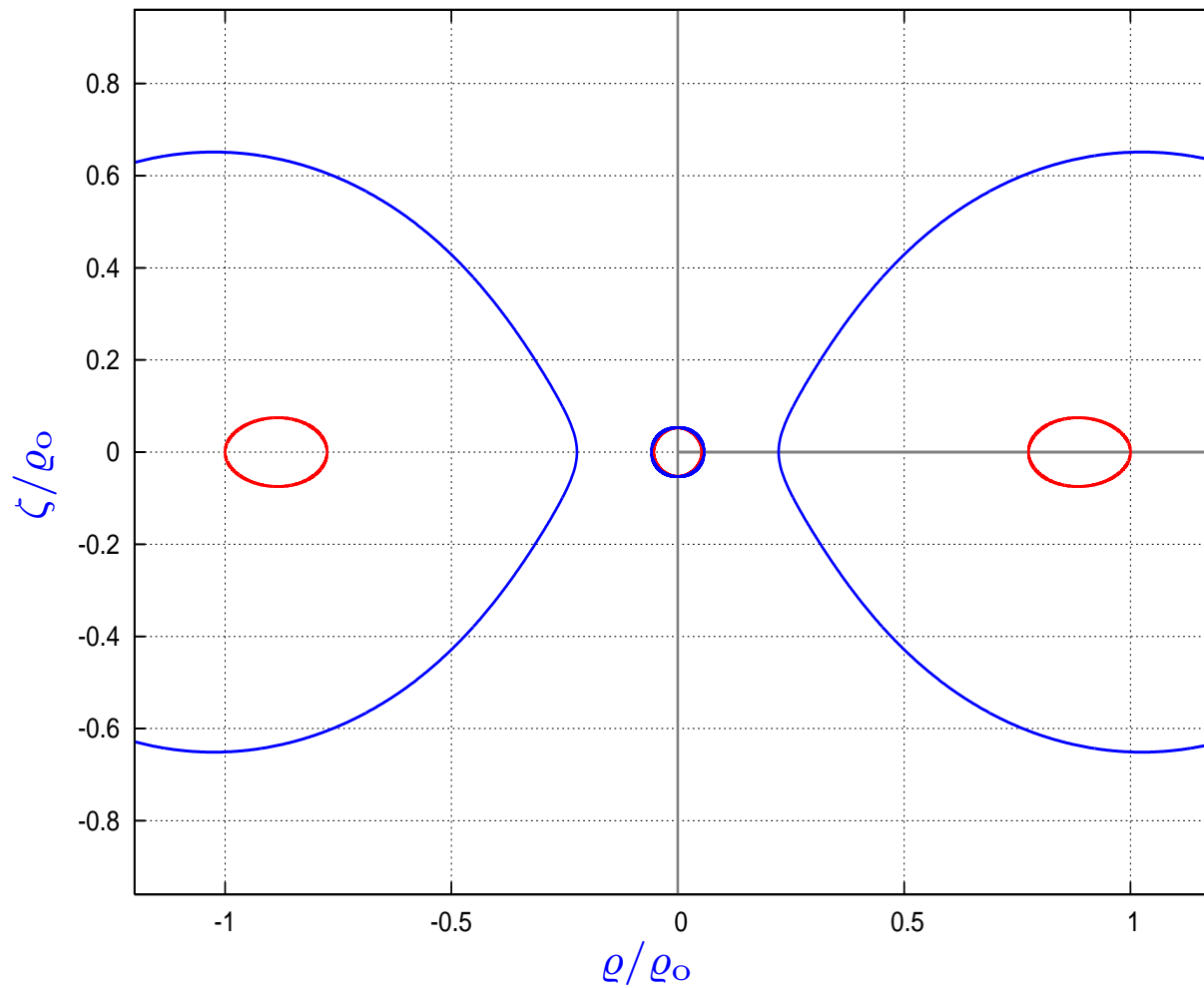
3. Journey

$$M_h/M_r = 0.094, Z_r = 3.6$$



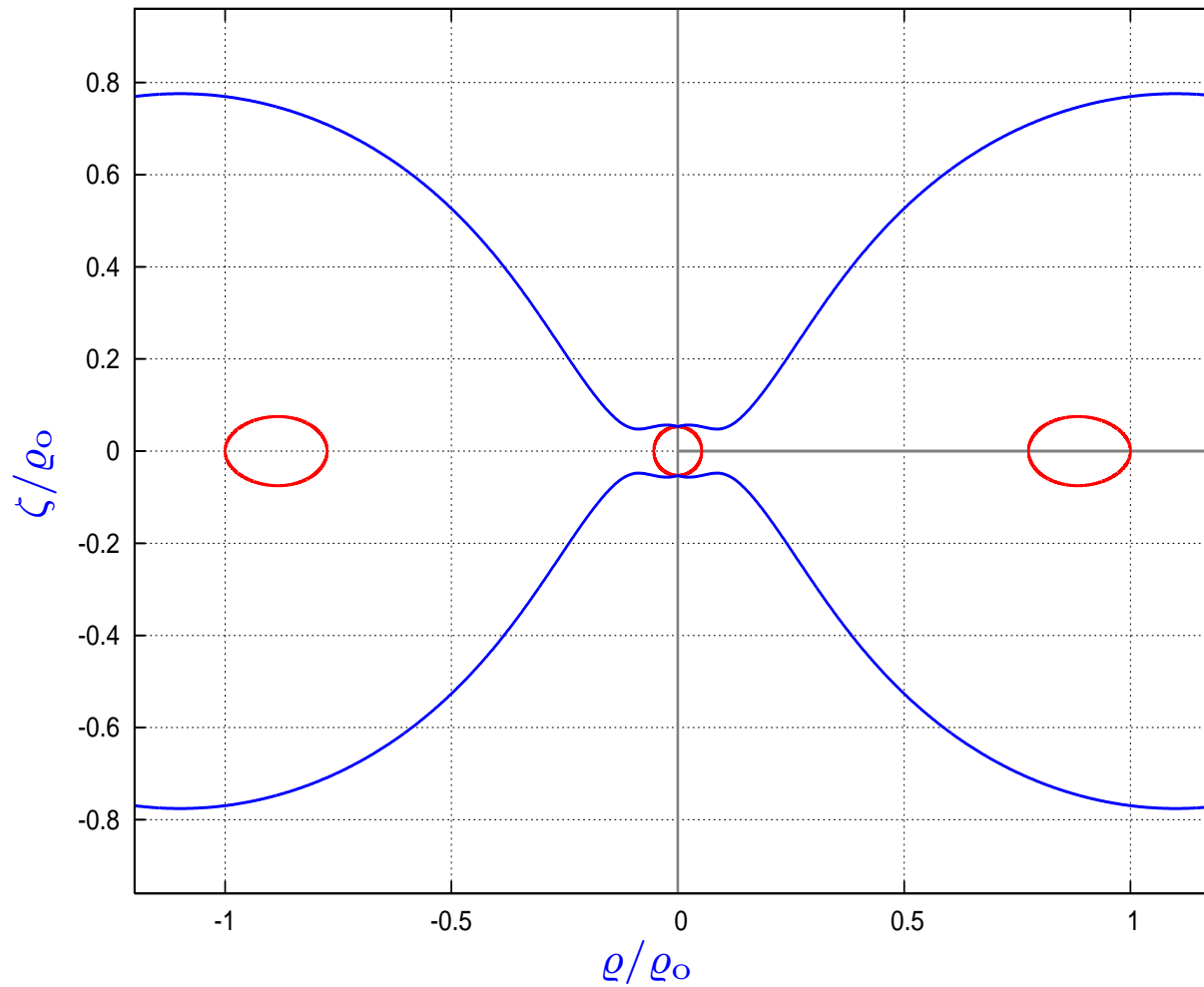
3. Journey

$$M_h/M_r = 0.069, Z_r = 4.2$$



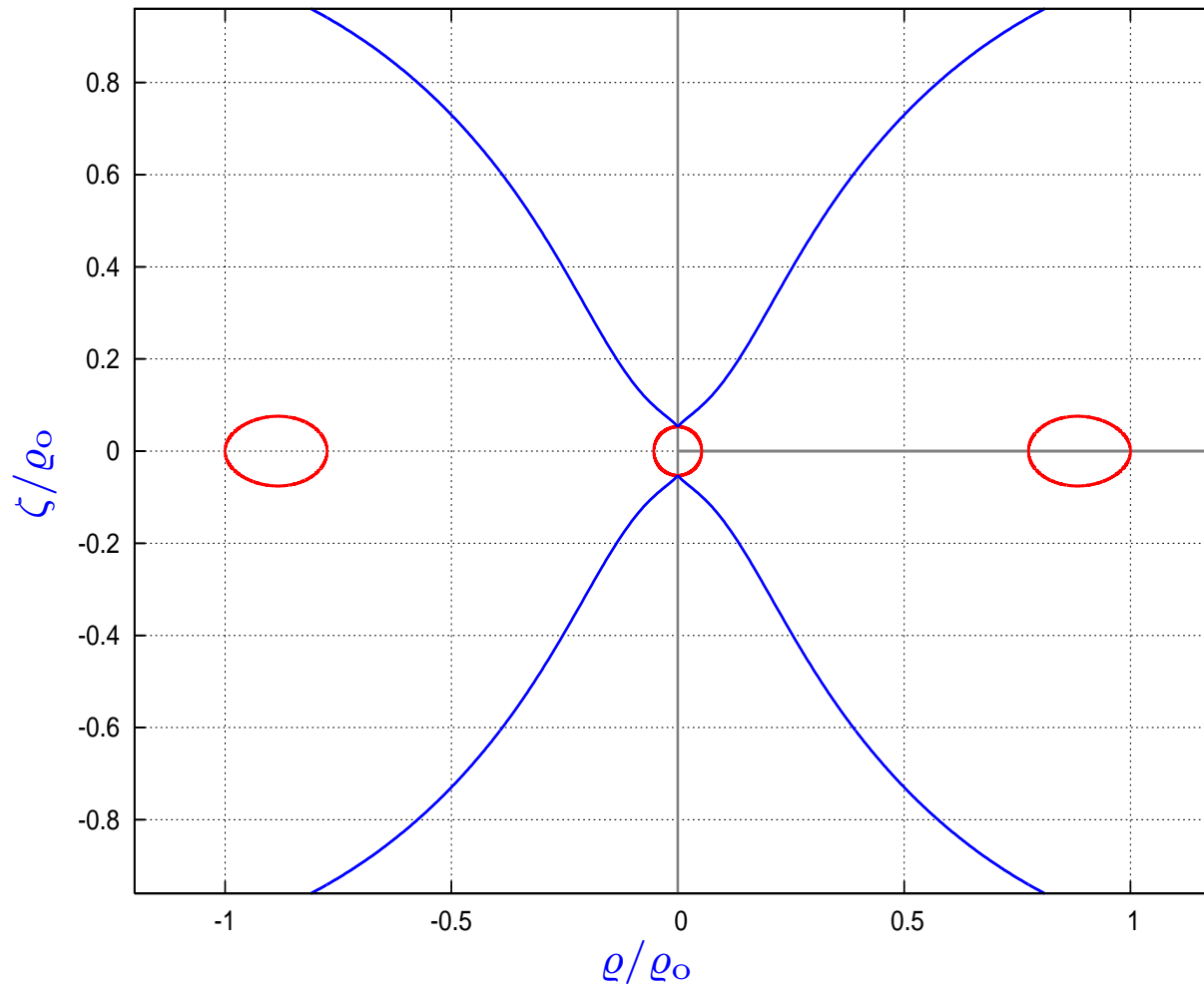
3. Journey

$$M_h/M_r = 0.048, Z_r = 4.8$$



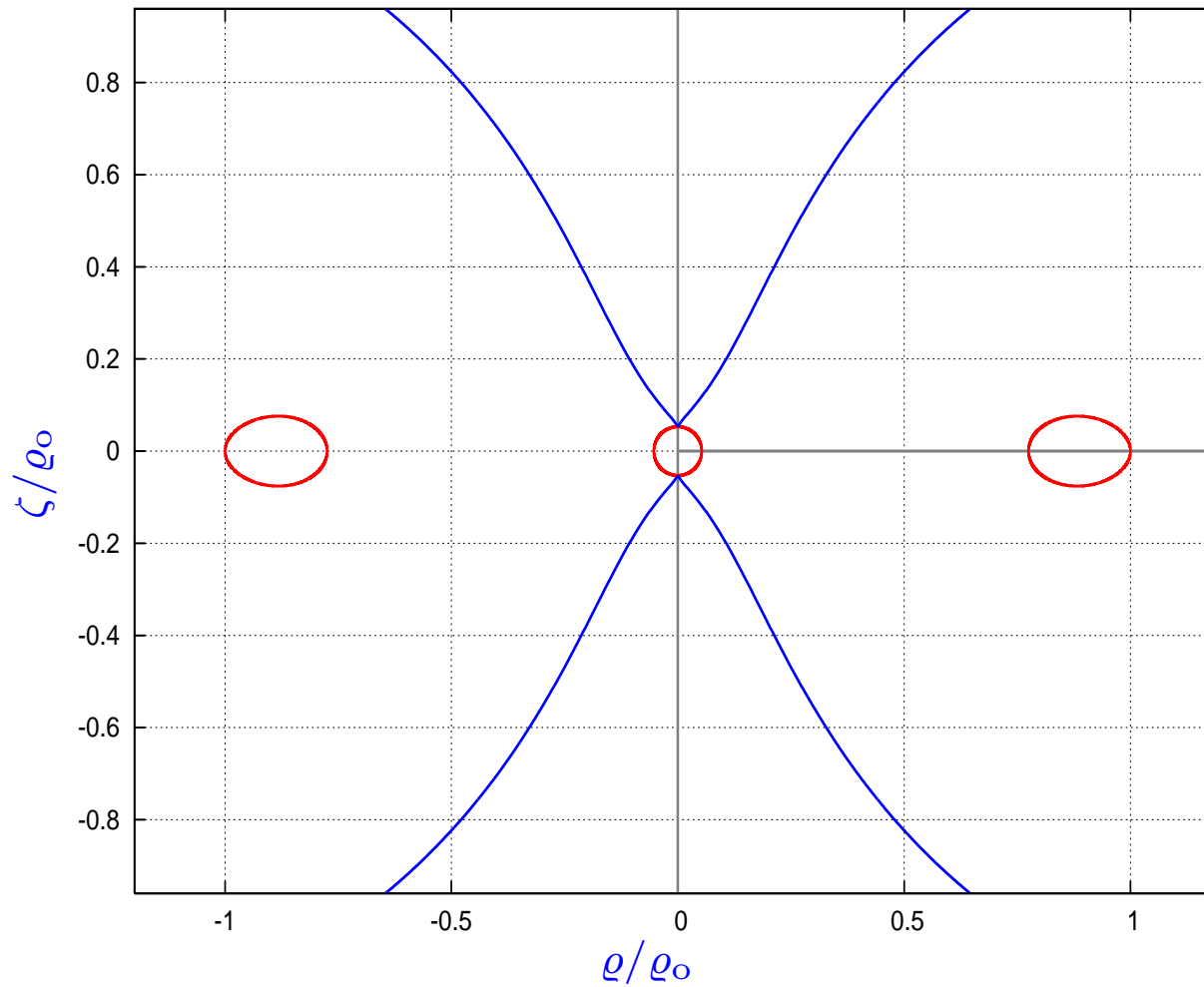
3. Journey

$$M_h/M_r = 0.013, Z_r = 6.4$$



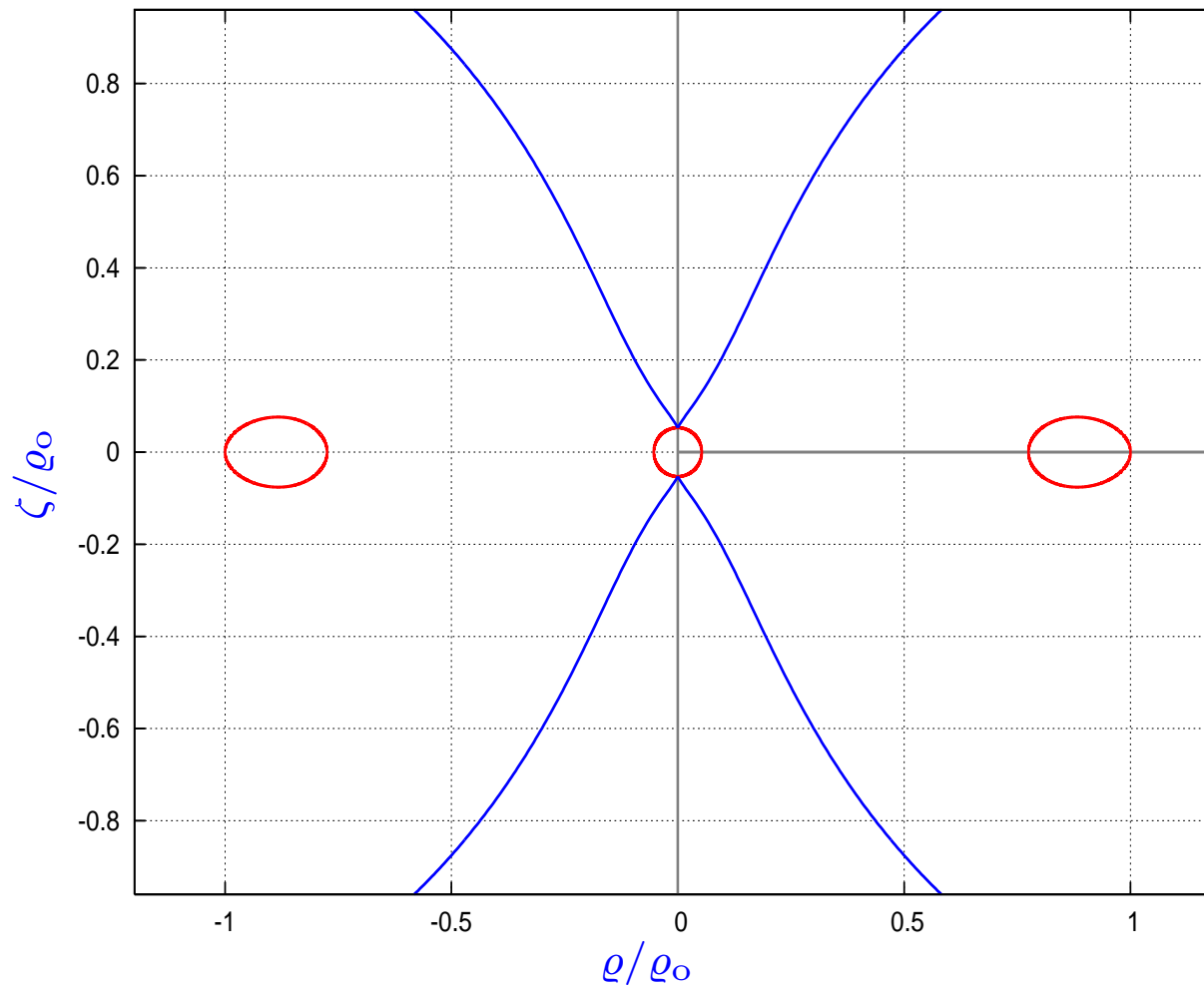
3. Journey

$$M_h/M_r = 0.00070, Z_r = 7.3$$



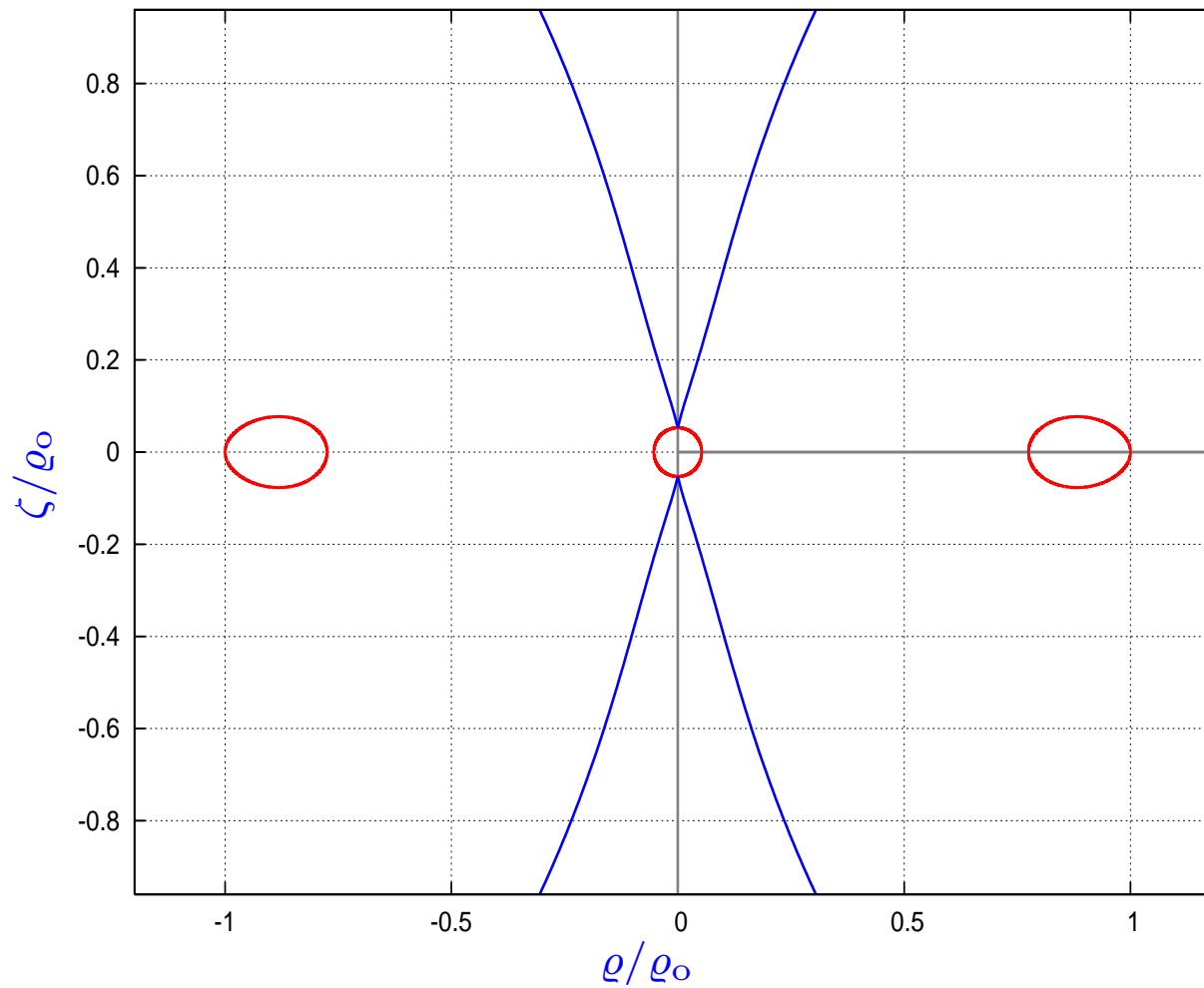
3. Journey

$$M_h/M_r = -0.0055, Z_r = 7.8$$



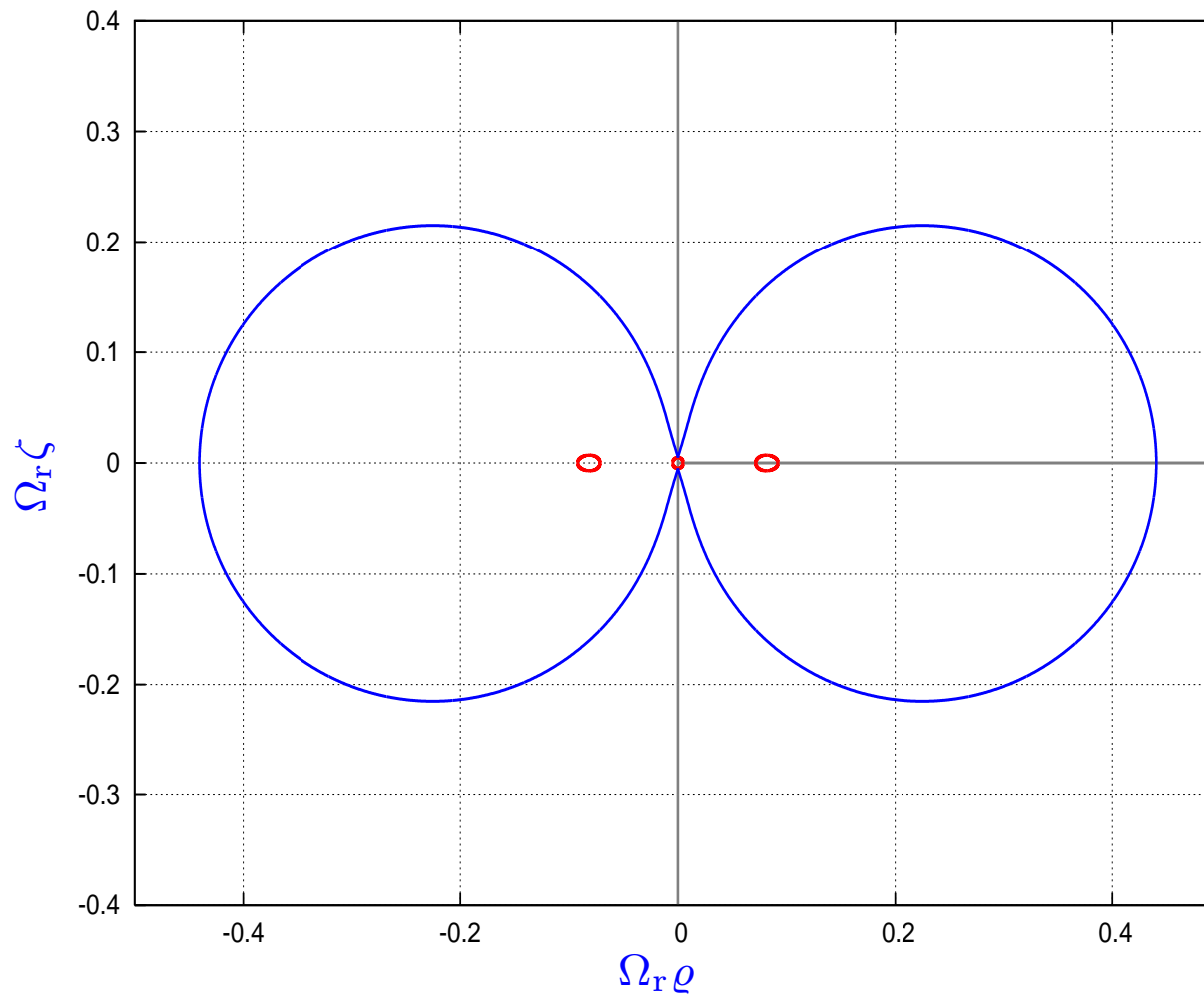
3. Journey

$$M_h/M_r = -0.04, Z_r = 13$$



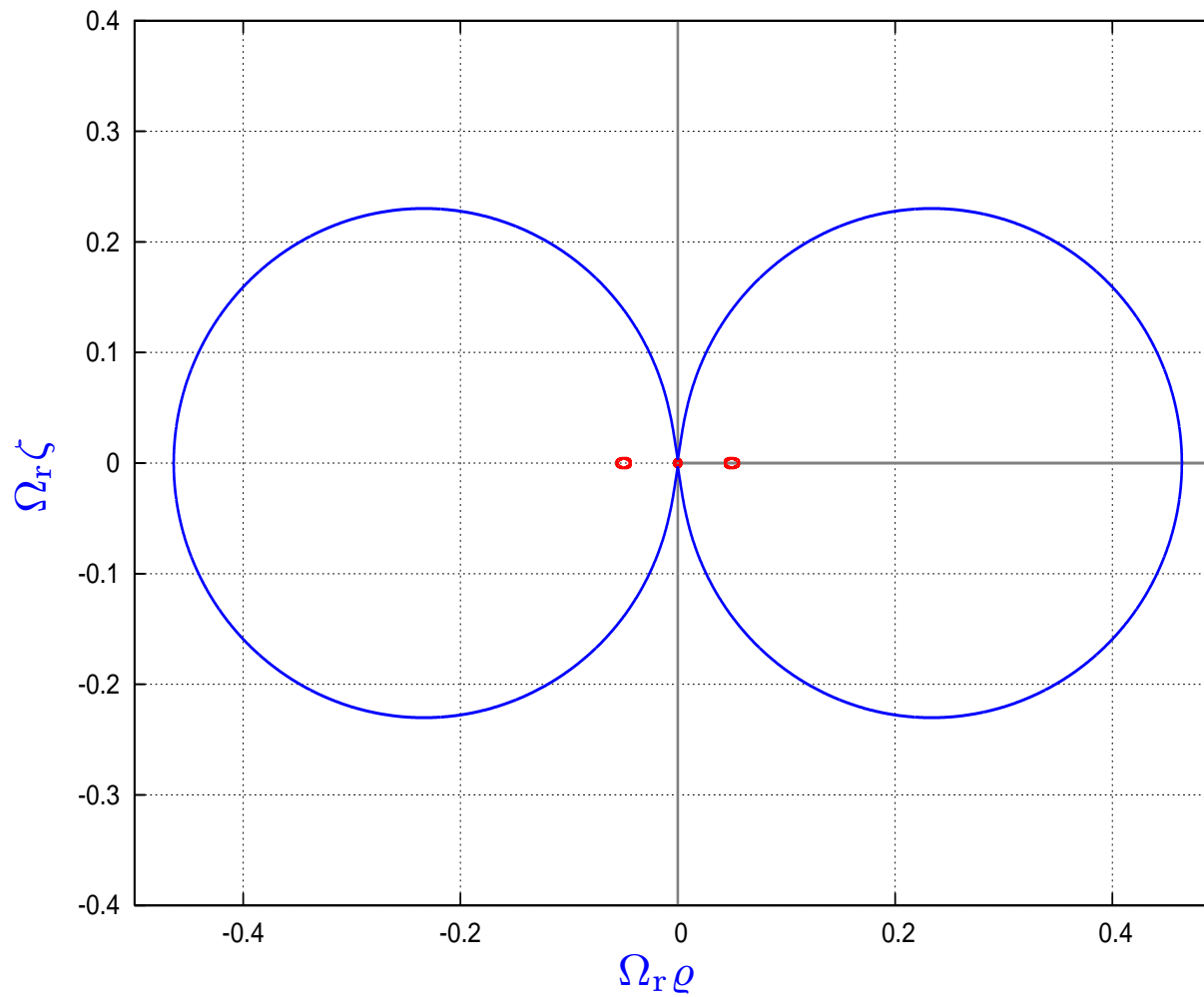
3. Journey

$$M_h/M_r = -0.04, Z_r = 13$$



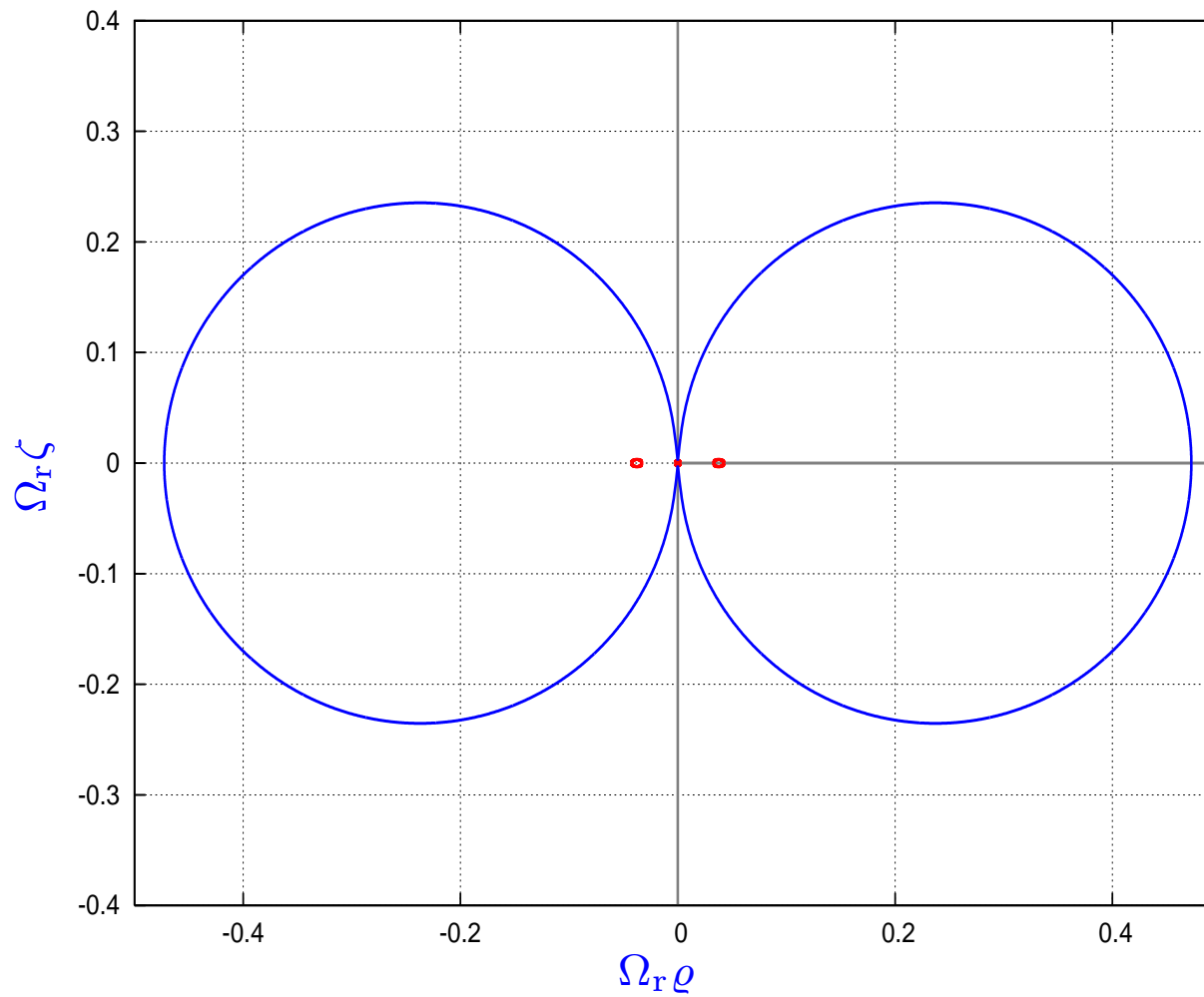
3. Journey

$$M_h/M_r = -0.060, Z_r = 24$$



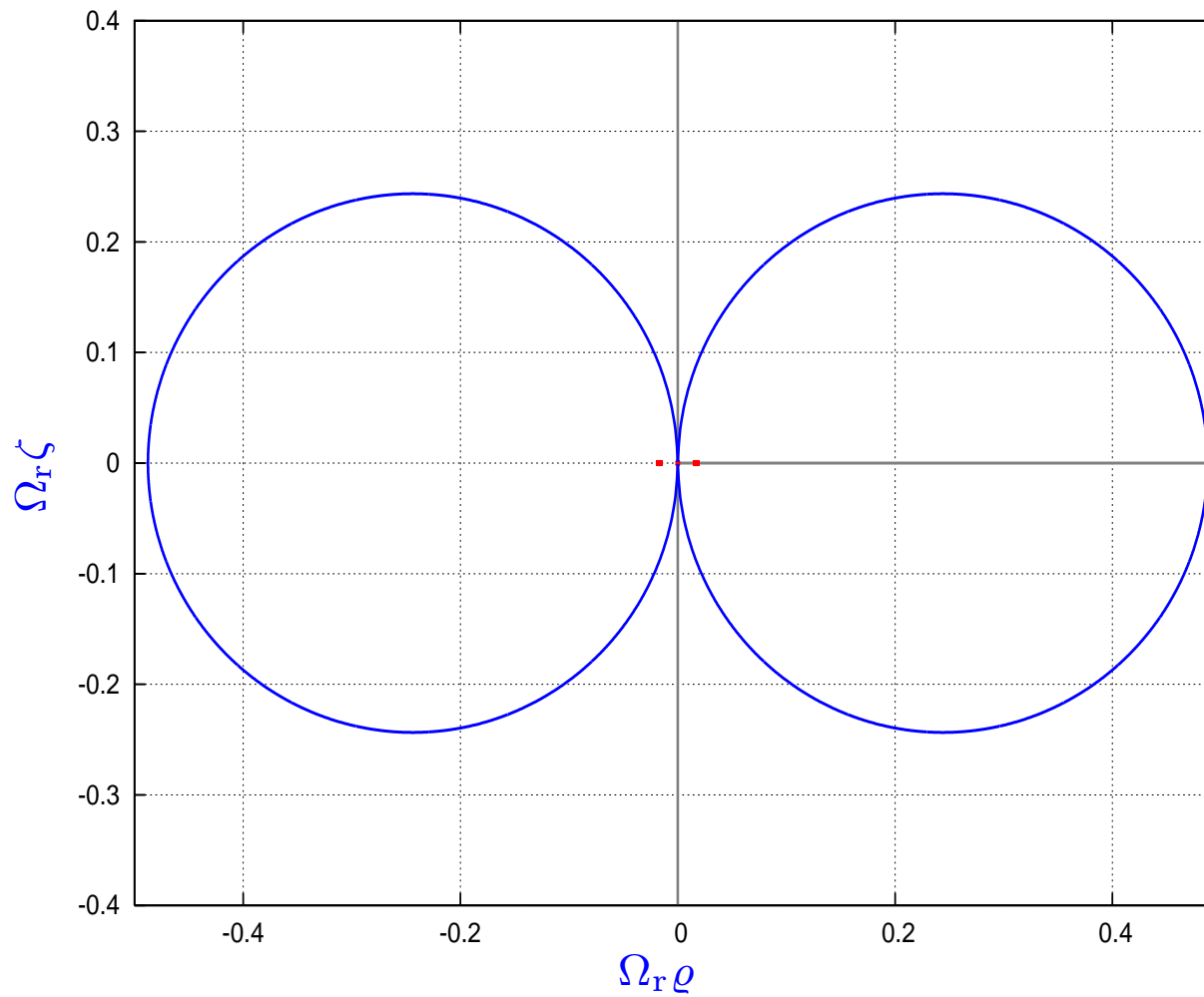
3. Journey

$$M_h/M_r = -0.067, Z_r = 32$$



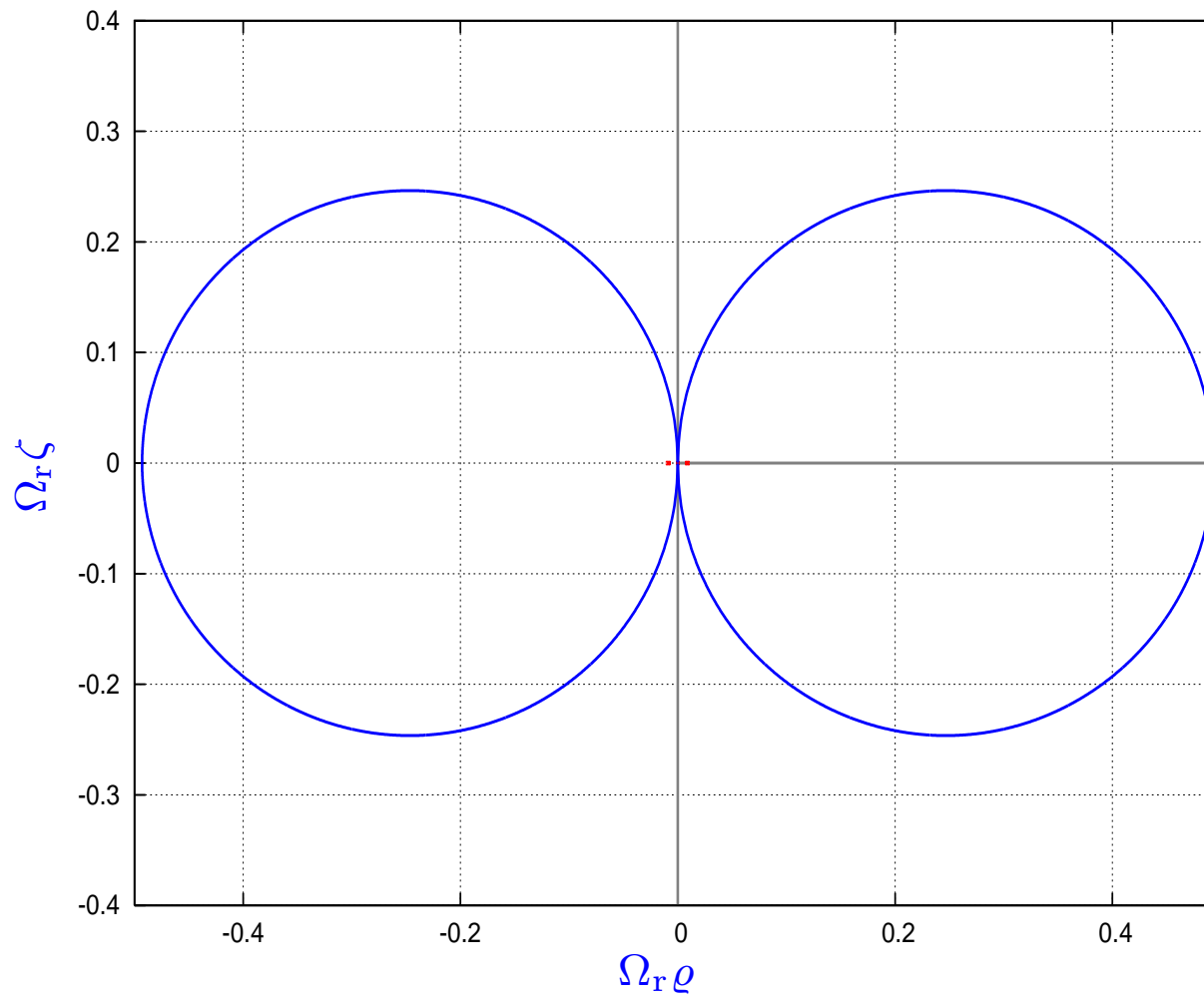
3. Journey

$$M_h/M_r = -0.077, Z_r = 75$$

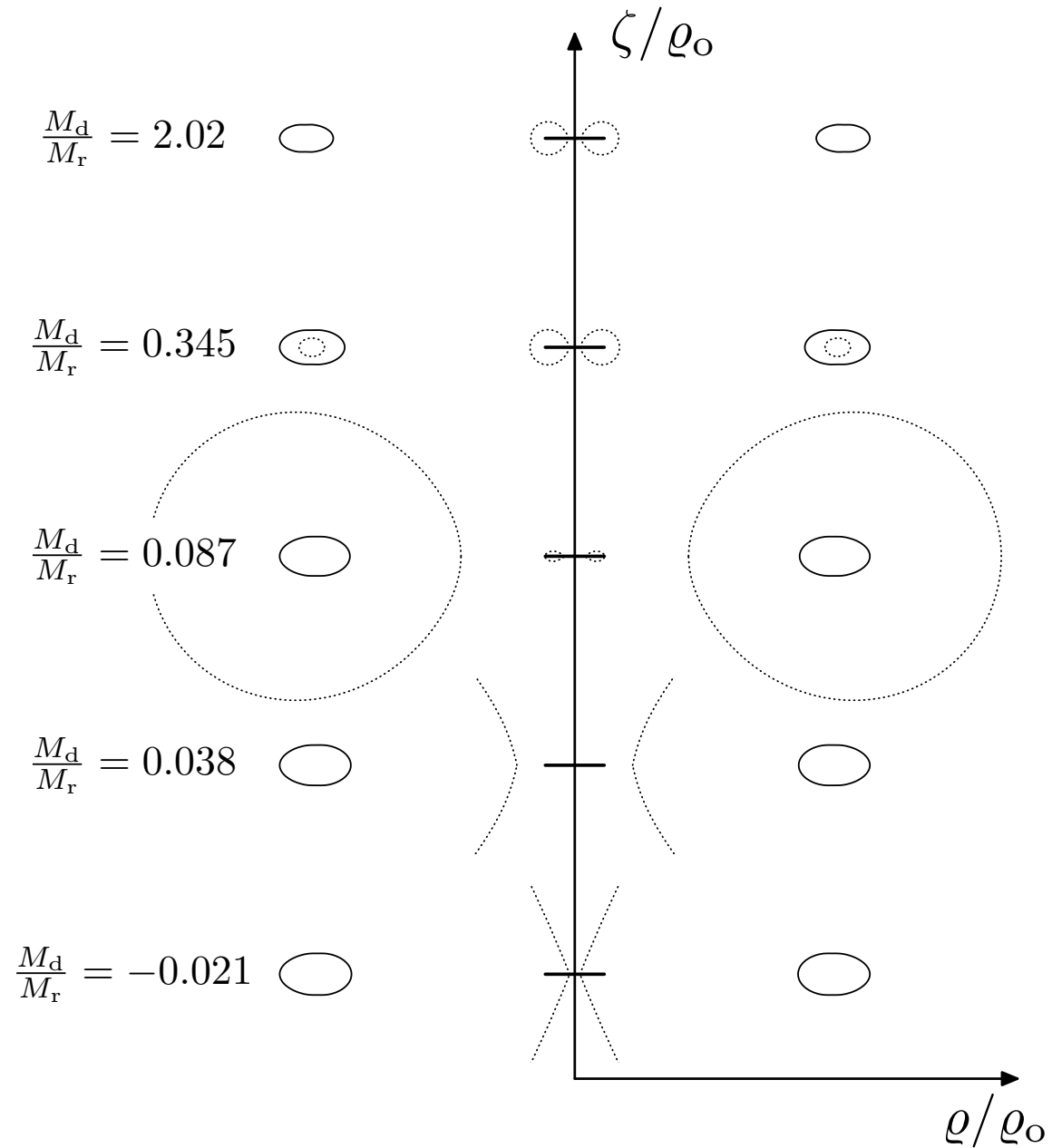


3. Journey

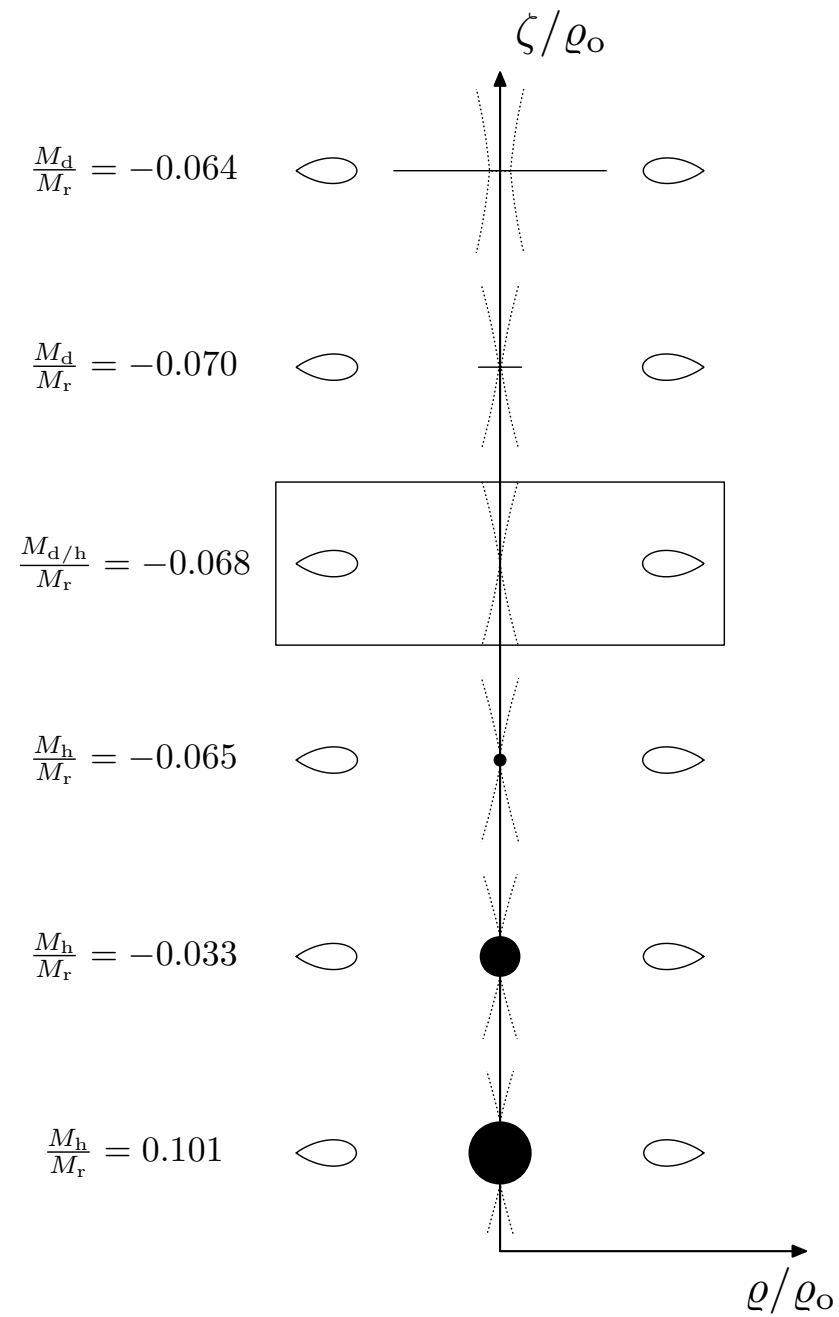
$$M_h/M_r = -0.080, Z_r = 150$$



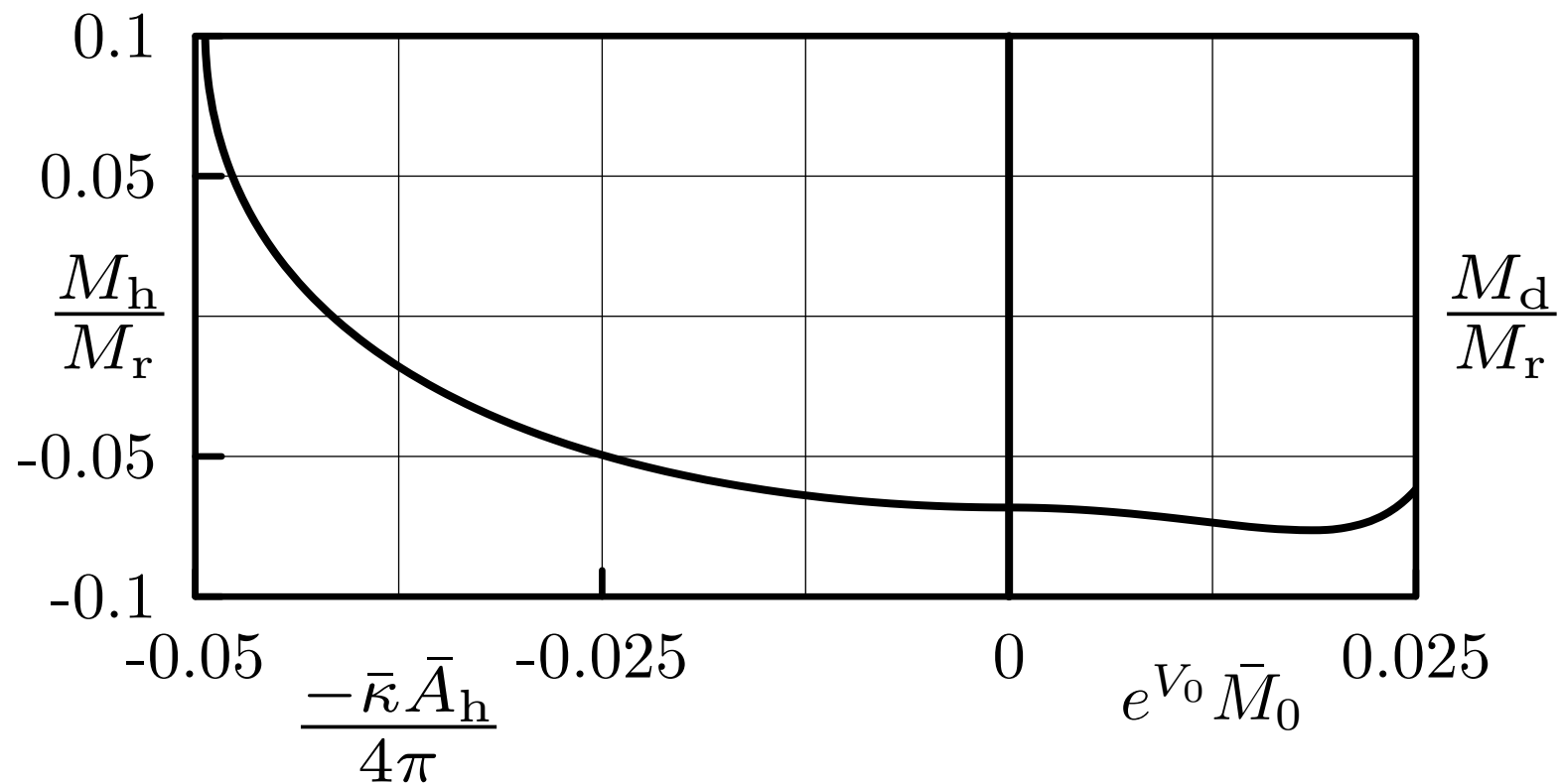
3. Journey



3. Journey



3. Journey



The ratio of the Komar Mass of the central object to that of the ring versus a measure of the distance to the degenerate black hole solution

4. Summary (1)

- The Komar mass of an object in axisymmetry and stationarity can be used on either side of the parametric transition from matter to a black hole.
- It can become negative if
 - (i) The object is placed within the strong gravitational field of a source with greater positive Komar mass.
 - (ii) This source is rapidly rotating so as to produce a large ergosphere encompassing the object.
 - (iii) The object is counter-rotating at a limited rate.
 - (iv) The object exerts a finite influence on the source (it is not close to a 'test'-object).

4. Summary (2)

- The Komar mass is **not** an intrinsic property of an object. It is a feature of an object within a specific highly relativistic spacetime geometry.
- **Question:** What is the maximally attainable ratio
$$-M^{\text{negative}} / M^{\text{positive}} \quad ?$$