

# Moving puncture method: recent results and outlook

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# Moving punctures work



# Methods

- Usual conformal decomposition:  $g_{ij} = \psi^A \gamma_{ij} = e^{4\phi} \gamma_{ij}$
- Initial Bowen-York puncture data:  
 $\psi = 1 + 0.5m_1/r_1 + 0.5m_2/r_2 + u$
- BSSN evolution:  $\gamma_{ij}$ ,  $\phi$ ,  $\Gamma^i$ ,  $A_{ij}$ ,  $K$
- Gauge: 1+log slicing, Gamma-driver shift allowed to be non-vanishing at the puncture
- 4<sup>th</sup> order Runge-Kutta, 4<sup>th</sup> order spatial differencing\*, 5<sup>th</sup> order interpolation at interfaces.
- Adaptive mesh refinement: finest around bh's.

# Gauge

$$\begin{aligned}\partial_t \alpha &= -2\alpha K + \beta^j \partial_j \alpha \\ \partial_t \beta^i &= \frac{3}{4} \tilde{\Gamma}^i + \beta^j \partial_j \beta^i - \eta \beta^i\end{aligned}$$

Note that if  $\eta = 0$ , the homogeneous equation,

$$\partial_t \beta^x - \beta^x \partial_x \beta^x = 0$$

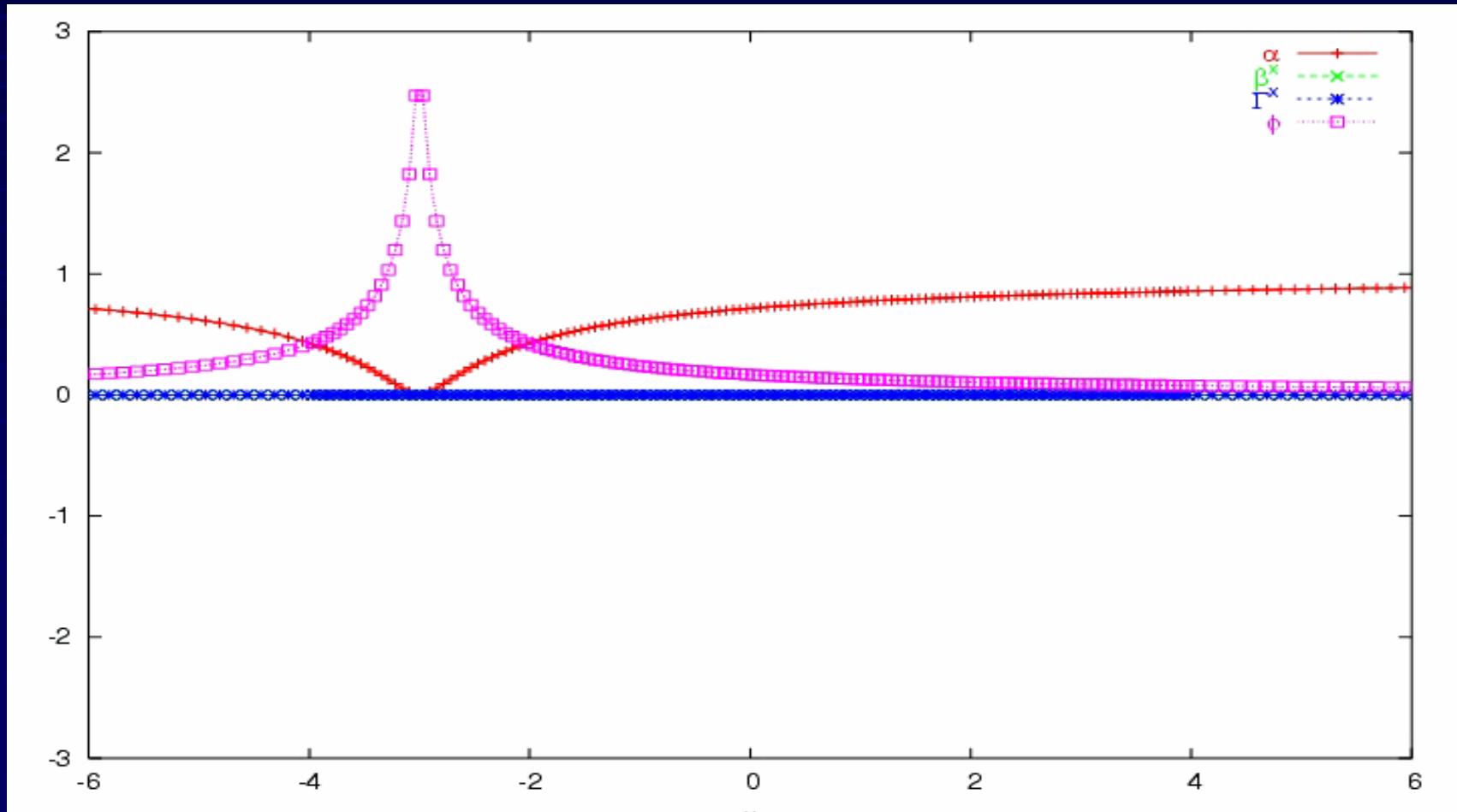
admits the shock wave solution,

$$\beta^x = \frac{x - x_0}{t_0 - t}$$

However, the source term  $\frac{3}{4} \tilde{\Gamma}^i$  is driven by derivatives of  $\beta^i$ :

$$\begin{aligned}\partial_t \tilde{\Gamma}^i &= \dots + \tilde{\gamma}^{jk} \partial_j \partial_k \beta^i + \frac{1}{3} \tilde{\gamma}^{ij} \partial_j \partial_k \beta^k \\ &\quad - \tilde{\Gamma}^j \partial_j \beta^i + \frac{2}{3} \tilde{\Gamma}^i \partial_j \beta^j + \dots\end{aligned}$$

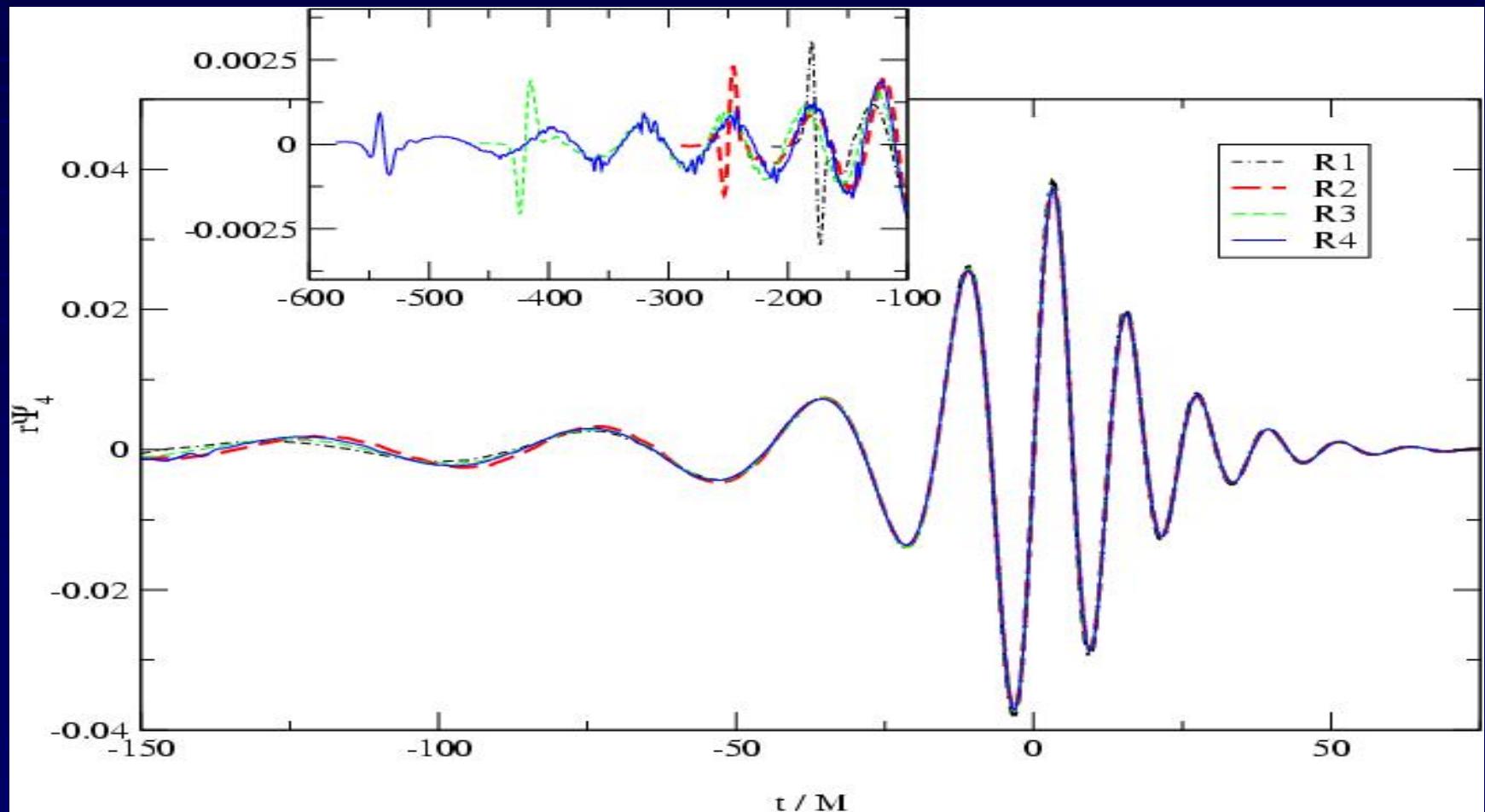
# Gauge test



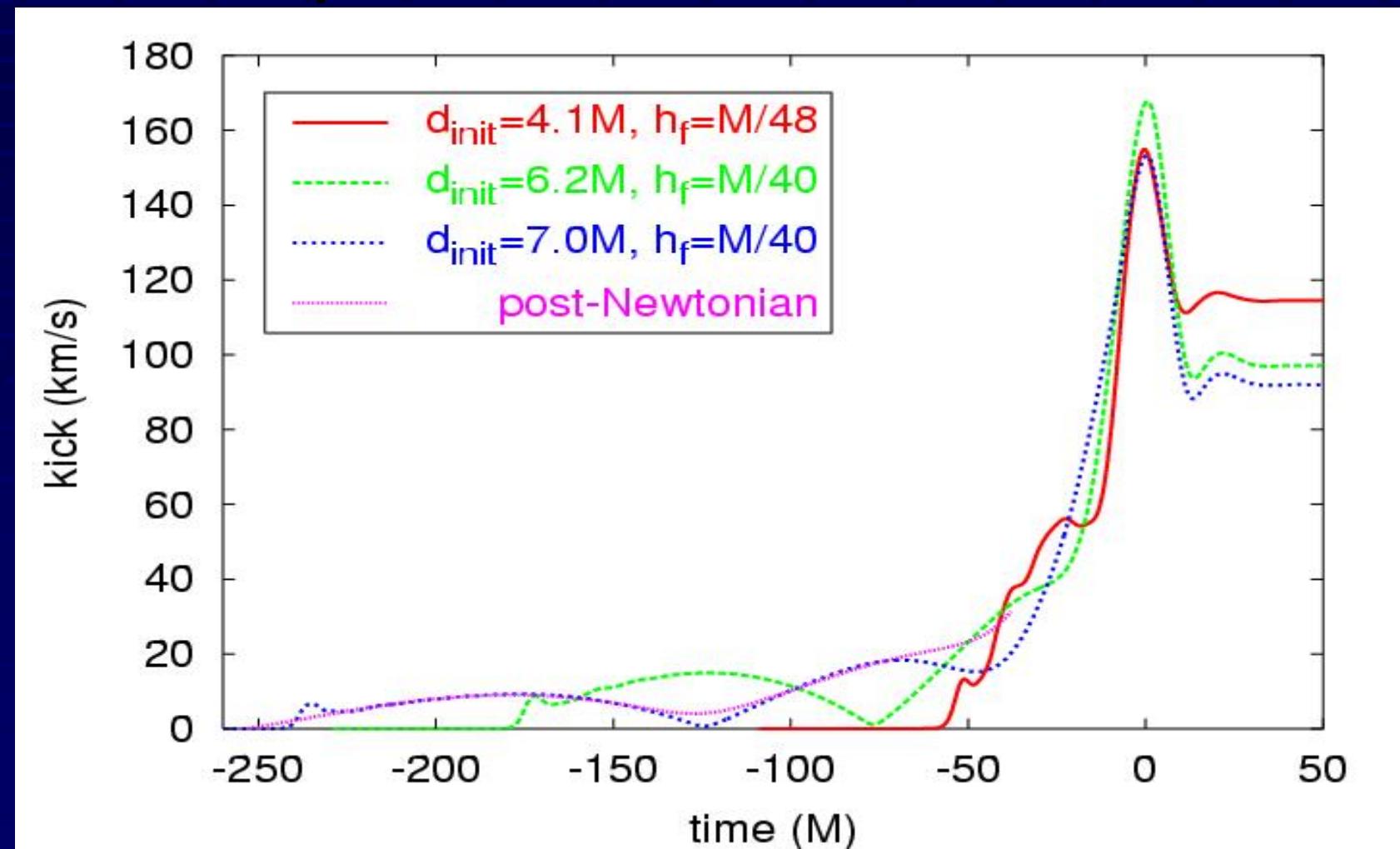
# Applications

- Equal-mass, various separations
  - Baker et al, Bruegmann et al, Campanelli et al
- Non-equal mass, recoil computation
  - Baker et al, Gonzalez et al, Herrmann et al
- Spins, aligned and anti-aligned
  - Campanelli et al

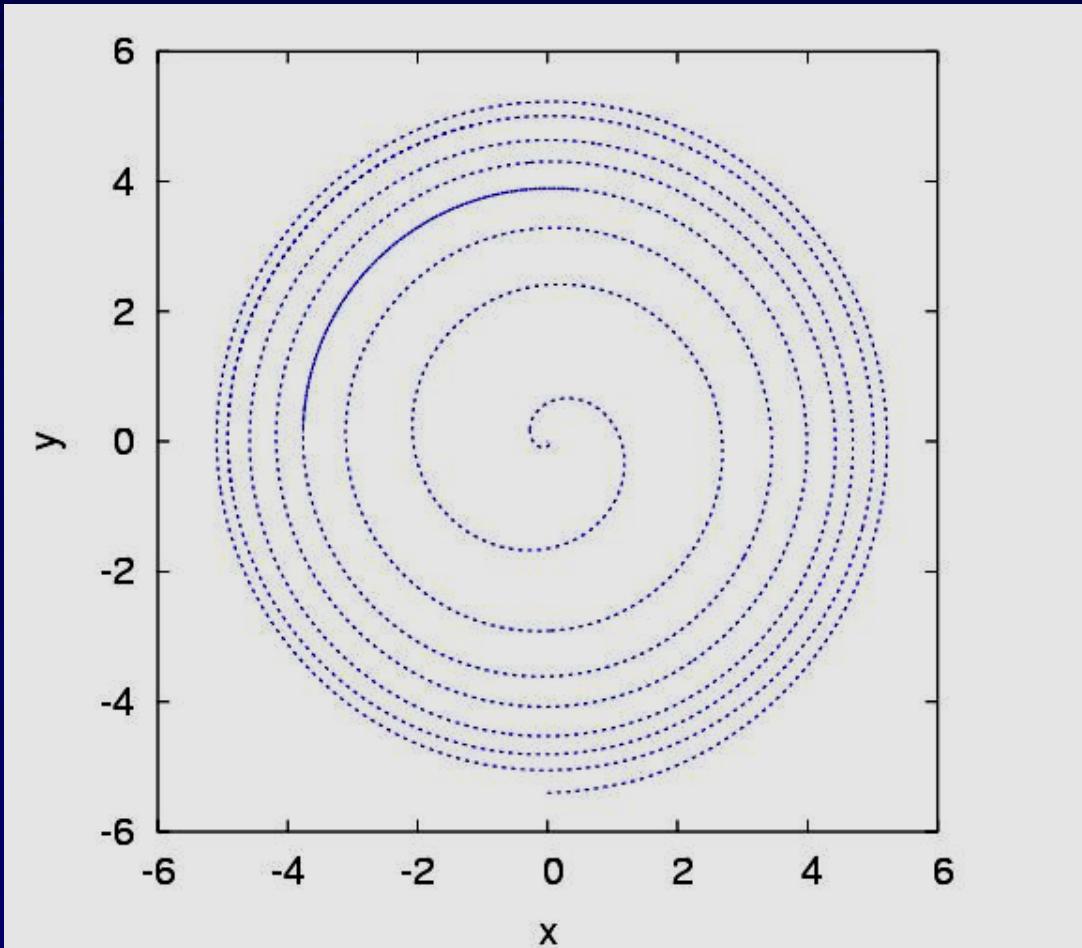
# Equal-mass, various separations



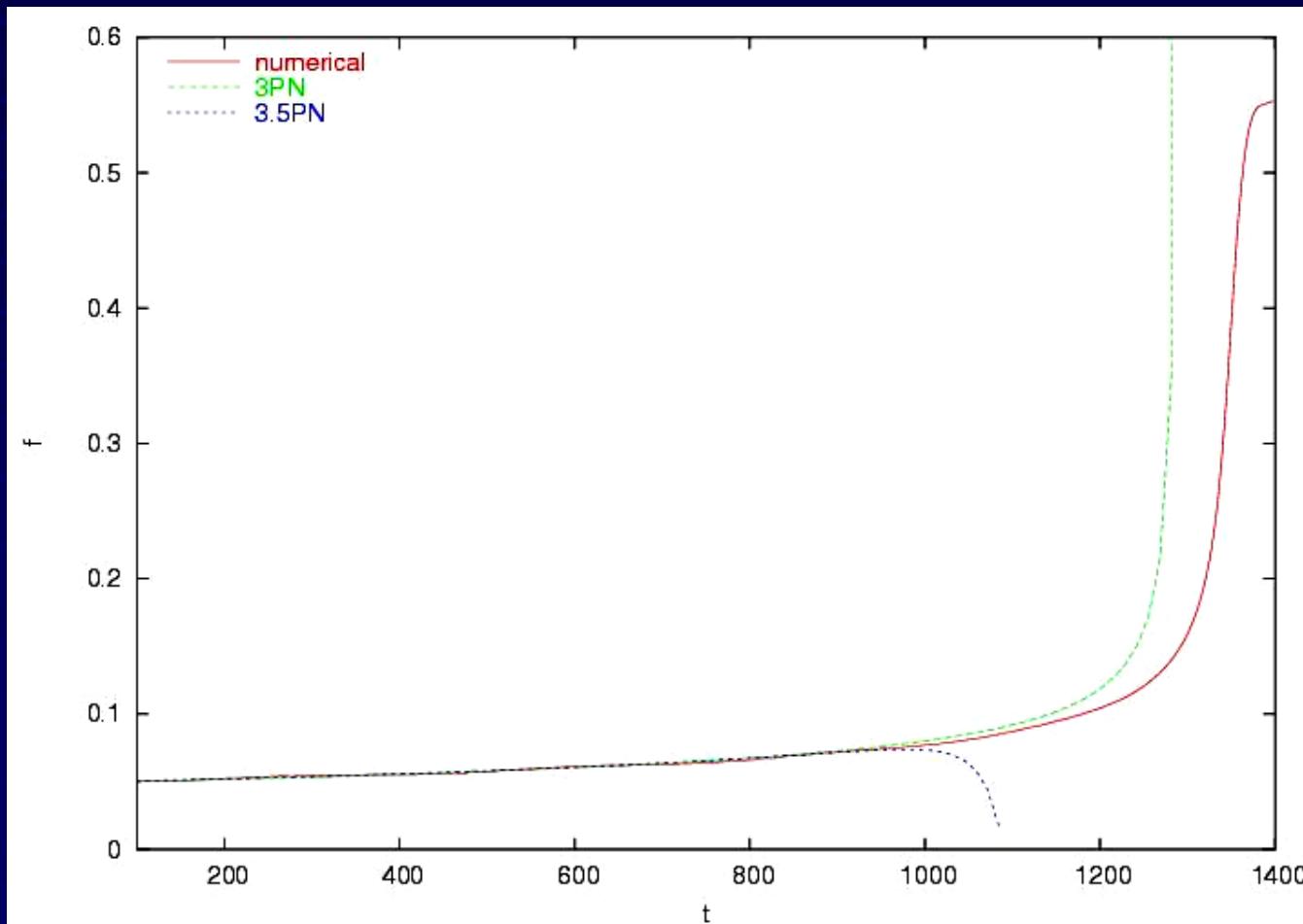
# Non-equal-mass, radiative recoil



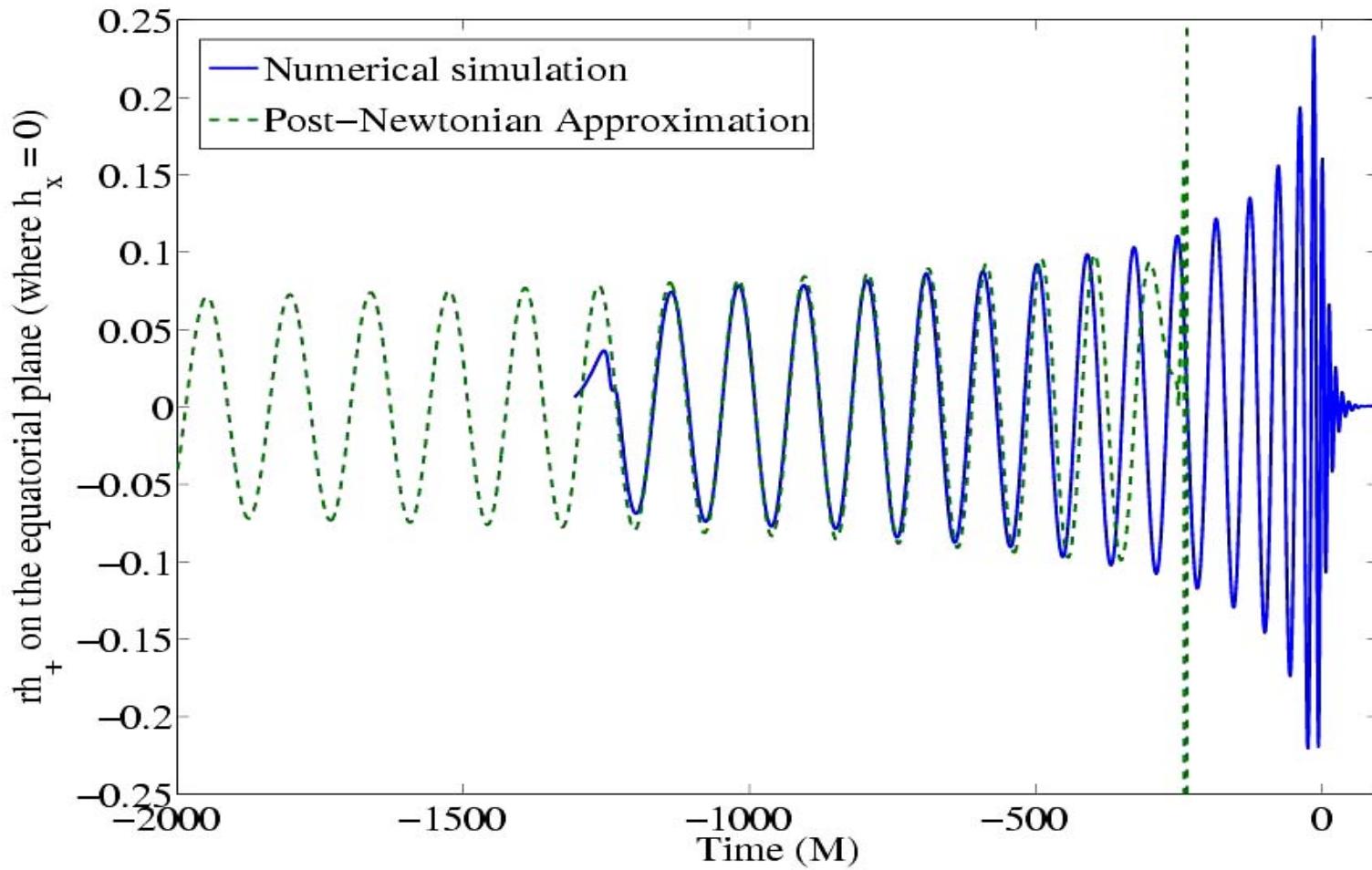
# Equal mass, 8-orbit run



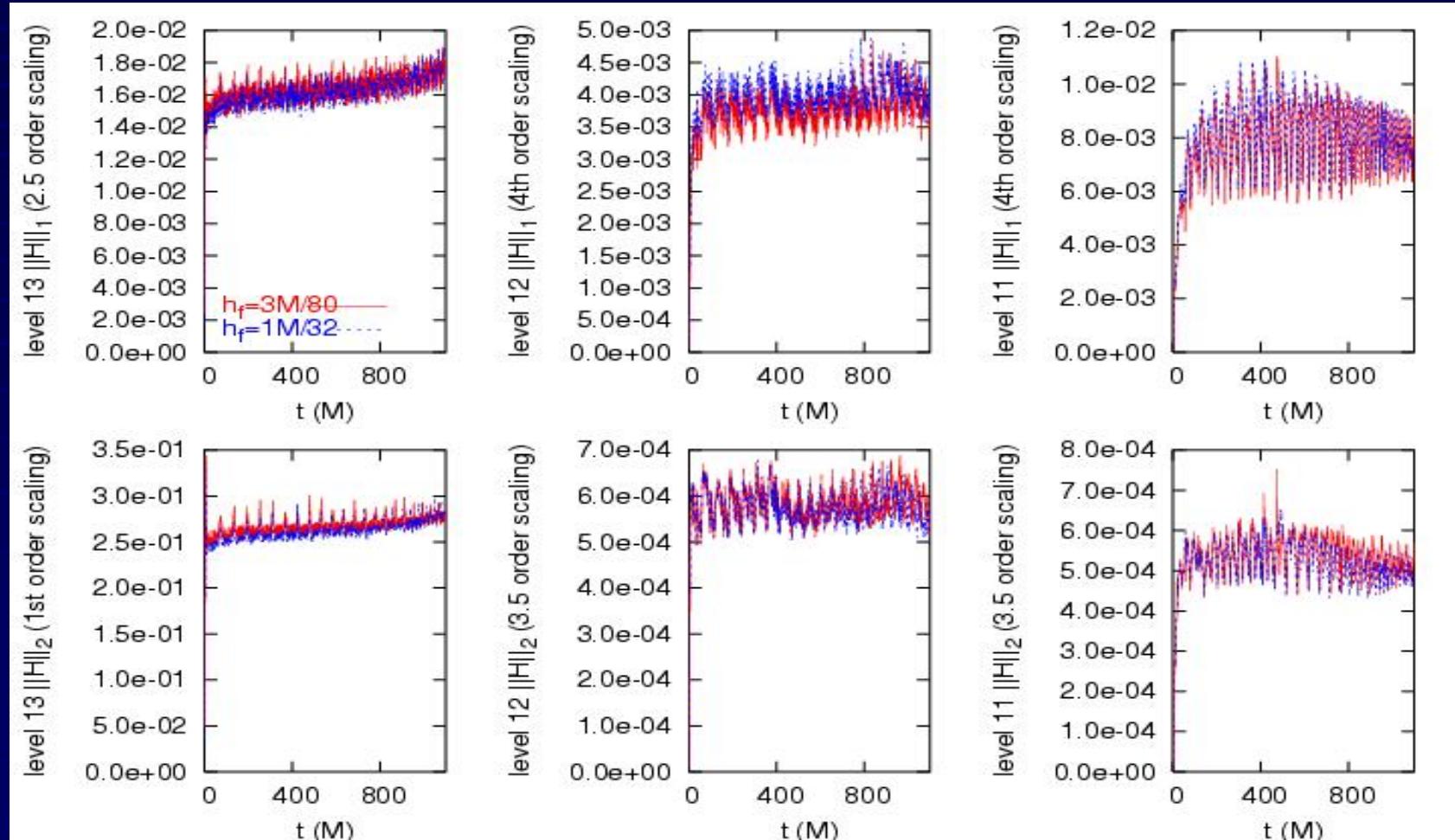
# GW frequency compared with PN



# $h+$ compared with PN



# Hamiltonian convergence



# Note on L1 and L2 norms

$$\begin{aligned}\|\text{error}\|_1 &\sim \frac{N^0 h^p + N^2 h^r + N^3 h^b}{N^3} \\ \|\text{error}\|_2 &\sim \sqrt{\frac{N^0 h^{2p} + N^2 h^{2r} + N^3 h^{2b}}{N^3}}\end{aligned}$$

$$N \sim h^{-1}$$

$p$  = order of puncture error

$r$  = order of refinement boundary error

$b$  = order of bulk error

Away from the puncture, if  $b = 4$  and  $r = 3$  then

$$\begin{aligned}\|H\|_1 &\sim h^4 \\ \|H\|_2 &\sim h^{3.5}\end{aligned}$$

But near the puncture, if  $p < 1$ , then

$$\begin{aligned}\|H\|_1 &\sim h^{p+3} \\ \|H\|_2 &\sim h^{\frac{2p+3}{2}}\end{aligned}$$

$$\psi \sim |\vec{x} - \vec{x}_{\text{punc}}|^{-\frac{1}{2}} \rightarrow p = -\frac{1}{2} ?$$

# Notes on puncture differencing

- So far, punctures have been confined to a plane between grid points.
- If error doesn't propagate from the puncture (e.g. via gauge mode) then it's not a concern.
- For differencing error within distance of  $O(h)$  from the puncture to be  $O(h^4)$ , fields must be  $C^4$ .
- Regarding the conformal factor variable:

If  $e^\phi \sim |\vec{x} - \vec{x}_{\text{punc}}|^{-\frac{1}{2}}$  then for positive, even  $n$ ,  $\chi \equiv e^{-n\phi}$  is  $C^{\frac{n}{2}-1}$  if  $n$  is not divisible by 4 and  $C^\infty$  if  $n$  is divisible by 4.

But note terms in the evolution equations of the form

$$e^{-4\phi} \partial_i \partial_j \phi = \frac{1}{n} \chi^{\frac{4}{n}} (\partial_i \chi \partial_j \chi / \chi^2 - \partial_i \partial_j \chi / \chi)$$

# A modified BSSN formulation

Assuming  $\lim_{\vec{x} \rightarrow \vec{x}_{\text{puncture}}} \alpha = O(|\vec{x} - \vec{x}_{\text{puncture}}|^2)$ , define:

$$\begin{aligned} w &= e^{-2\phi} \\ q &= e^{2\phi}\alpha \end{aligned}$$

where:

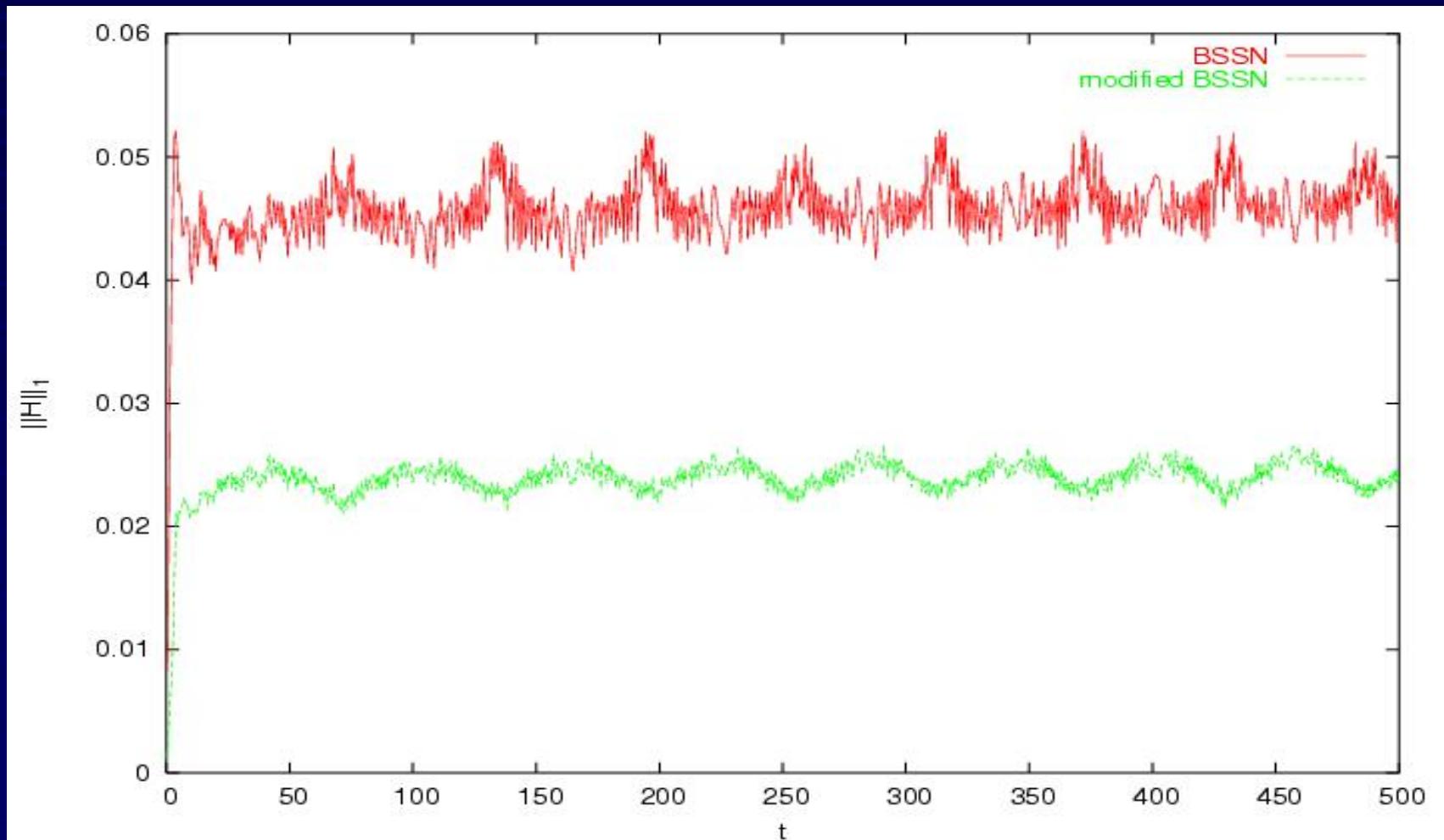
$$\begin{aligned} w \nabla_m \nabla_n (qw) &= w \partial_m \partial_n (qw) - 4w \partial_{(m} \phi \partial_{n)} (qw) \\ &\quad - w \bar{\Gamma}_{mn}^k \partial_k (qw) + 2\bar{\gamma}_{mn} \bar{\gamma}^{kl} w \partial_k \phi \partial_l (qw) \end{aligned}$$

Then the BSSN equations become:

$$\begin{aligned} \partial_t w &= \frac{1}{3}qw^2 K - \frac{1}{3}w \partial_k \beta^k + \beta^k \partial_k w \\ \partial_0 K &= -\bar{\gamma}^{ij} w^2 \nabla_j \nabla_i (qw) + qw \left( \bar{A}_{ab} \bar{A}^{ab} + \frac{1}{3}K^2 \right) & w^2 R_{ij} &= w^2 \bar{R}_{ij} - 2w^2 \bar{\nabla}_i \bar{\nabla}_j \phi - 2\bar{\gamma}_{ij} w^2 \bar{\nabla}^k \bar{\nabla}_k \phi \\ &\quad + 4w \partial_i \phi w \partial_j \phi - 4\bar{\gamma}_{ij} w \partial^k \phi w \partial_k \phi \\ \partial_0 \bar{\gamma}_{ij} &= -2qw \bar{A}_{ij} & w^2 \bar{\nabla}_i \bar{\nabla}_j \phi &= w^2 \partial_i \partial_j \phi - \bar{\Gamma}_{ij}^k w^2 \partial_k \phi \\ \partial_0 \bar{A}_{ij} &= w \left[ -w \nabla_i \nabla_j (qw) + qw^2 R_{ij} \right]^{\text{TF}} \\ &\quad + qw \left( K \bar{A}_{ij} - 2\bar{A}_{ia} \bar{A}^a{}_j \right) & w^2 \partial_i \partial_j \phi &= \frac{1}{2} (\partial_i w \partial_j w - w \partial_i \partial_j w) \\ \partial_t \bar{\Gamma}^i &= 2q \left( w \bar{\Gamma}_{ab}^i \bar{A}^{ab} - \frac{2}{3}w \bar{\gamma}^{ia} K_{,a} + 6\bar{A}^{ia} w \phi_{,a} \right) \\ &\quad + \bar{\gamma}^{kl} \left( -\bar{\Gamma}_{kl}^j \beta^i{}_j + \frac{2}{3} \bar{\Gamma}_{kl}^i \beta^j{}_j \right) + \beta^k \bar{\Gamma}^i{}_{,k} \\ &\quad + \bar{\gamma}^{jk} \beta^i{}_{jk} + \frac{1}{3} \bar{\gamma}^{ij} \beta^k{}_{kj} - 2\bar{A}^{ia} (qw)_{,a} & w \partial_i \phi &= -\frac{1}{2} \partial_i w \\ \partial_t q &= -2qK - \frac{1}{3}q^2 wK + \frac{1}{3}q \partial_k \beta^k + \beta^k \partial_k q & & \end{aligned}$$

$$\partial_t \beta^i = \frac{3}{4} \bar{\Gamma}^i + \beta^j \partial_j \beta^i - \eta \beta^i$$

# Modified BSSN test



# Recap and open questions

- Moving punctures yield stable and accurate evolutions in a variety of cases.
- Does error ever propagate from punctures?
- What is the best lapse condition?
  - Asymptotic behavior near puncture
  - Gauge speed
- What is the best variant of BSSN?
  - $\chi = \psi^n$  ?
  - Densitize other variables to make  $C^4$  ?