

Detecting bodies orbiting the Galactic Center black hole Sgr A* with LISA

Éric Gourgoulhon¹, Alexandre Le Tiec¹, Frédéric Vincent², Niels Warburton³

¹ LUTH, Observatoire de Paris, Université PSL, CNRS, Univ. de Paris, 92190 Meudon, France

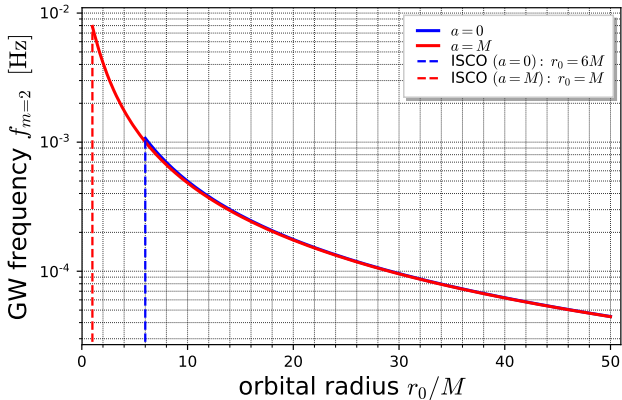
² LESIA, Observatoire de Paris, Université PSL, CNRS, Sorbonne Université, Univ. de Paris, 92190 Meudon, France

³ School of Mathematics and Statistics, University College Dublin, Belfield, Dublin 4, Ireland

based on [A&A 627, A92 \(2019\) \[arXiv:1903.02049\]](#)

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GW frequencies from circular orbits around Sgr A*



Angular velocity of circular equatorial orbits around a Kerr BH

$$\omega_0 = \frac{M^{1/2}}{r_0^{3/2} + aM^{1/2}}$$

Dominant GW frequency

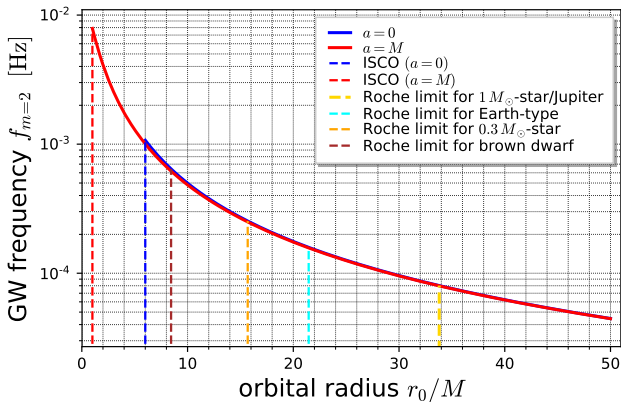
$$f_{m=2} = 2f_0 = \frac{\omega_0}{\pi}$$

Sgr A* mass

$$\begin{aligned} M &= 4.10 \times 10^6 M_\odot \\ &= 20.2 \text{ s} \end{aligned}$$

[Gravity team, *A&A* 615, L15 (2018)]

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Roche radius: $r_R \simeq 1.14 \left(\frac{M}{\rho} \right)^{1/3}$

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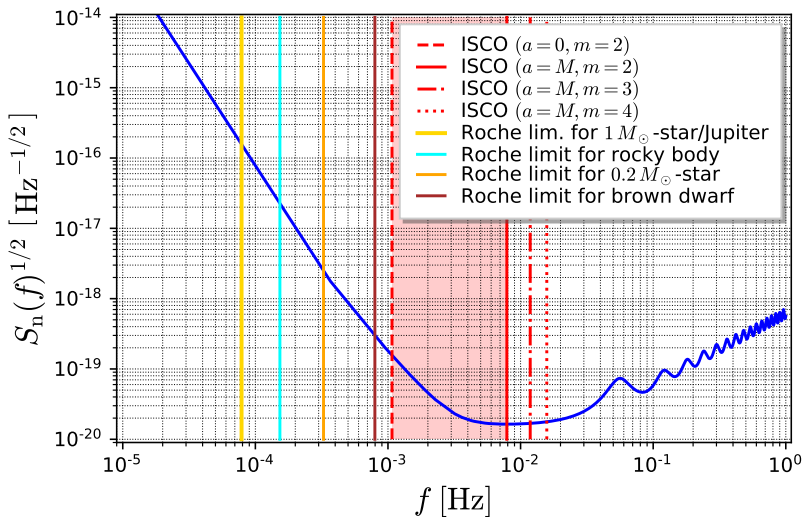
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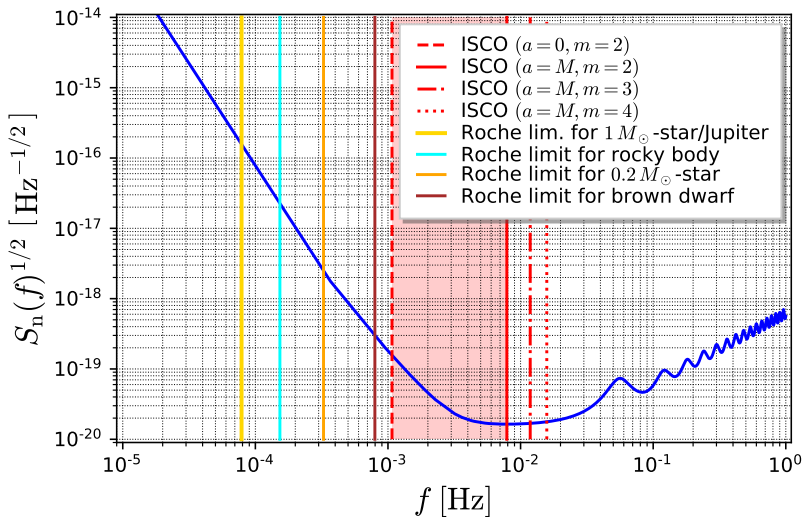
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Frequencies of Sgr A* close orbits are in LISA band



ISCO for $a = M$: $f_{m=2} = 7.9$ mHz

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ISCO for $a = M$: $f_{m=2} = 7.9 \text{ mHz}$ ← coincides with LISA max. sensitivity!

Previous studies of Sgr A* as a source for LISA

- Freitag (2003) [ApJ 583, L21]: GW from orbiting stars at quadrupole order; low-mass main-sequence (MS) stars are good candidates for LISA
- Barack & Cutler (2004) [PRD 69, 082005]: $0.06M_{\odot}$ MS star observed 10^6 yr before plunge \implies SNR = 11 in 2 yr of LISA data \implies Sgr A*'s spin within 0.5% accuracy
- Berry & Gair (2013) [MNRAS 429, 589]: extreme-mass-ratio burst (single periastron passage on a highly eccentric orbit) \implies GW burst \implies LISA detection of $10M_{\odot}$ for periastron $< 65M$; event rate could be $\sim 1 \text{ yr}^{-1}$
- Linial & Sari (2017) [MNRAS 469, 2441]: GW from orbiting MS stars undergoing Roche lobe overflow \implies detectability by LISA; possibility of a *reverse chirp signal (outspiral)*
- Amaro-Seoane (2019) [PRD 99, 123025]: *Extremely Large Mass-Ratio Inspirals (X-MRI)* \implies brown dwarfs orbiting Sgr A* should be detected in great numbers by LISA: ~ 20 in band at any time
- Kühnel et al. (2020) [EPJ C 80, 627]: GW from an ensemble of macroscopic dark matter candidates orbiting Sgr A*, such as primordial BHs, with masses in the range $10^{-13} - 10^3 M_{\odot}$

Our study

All previous studies have been performed in a Newtonian framework (quadrupole formula). Now, for orbits close to the ISCO, relativistic effects are expected to be important.

⇒ we have adopted a **fully relativistic framework**:

- Sgr A* is modeled as a Kerr BH and GW are computed via the theory of perturbations of the Kerr metric
- tidal effects are evaluated via the theory of Roche potential in the Kerr metric developed by Dai & Blandford (2013) [[MNRAS 434, 2948](#)]

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Limitation: **circular equatorial orbits**; valid for

- inspiralling compact objects arising from the tidal disruption of a binary (*zero-eccentricity EMRI*)
- main-sequence stars formed in an accretion disk
- compact objects resulting from the most massive of such stars
- $\sim 1/4$ of the population of brown dwarfs studied by Amaro-Seoane (2019)

Waveforms from circular orbits

computed as linear perturbations of Kerr metric (Teukolsky 1973)

Detweiler (1978)

$$h_+ - ih_\times = \frac{2\mu}{r} \sum_{\ell=2}^{\infty} \sum_{\substack{m=-\ell \\ m \neq 0}}^{\ell} \frac{Z_{\ell m}^{\infty}(r_0)}{(m\omega_0)^2} {}_{-2}S_{\ell m}^{am\omega_0}(\theta, \varphi) e^{-im(\omega_0(t-r_*)+\varphi_0)}$$

μ : mass of orbiting object; (t, r, θ, φ) : Boyer-Lindquist coordinates of the observer

${}_{-2}S_{\ell m}^{am\omega_0}(\theta, \varphi)$: spheroidal harmonics of spin weight -2

Waveforms from circular orbits

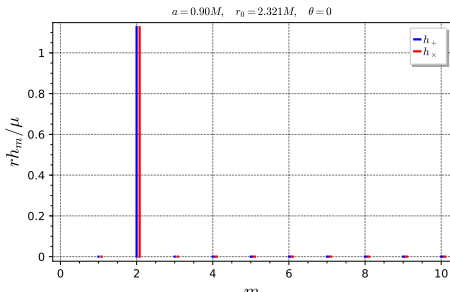
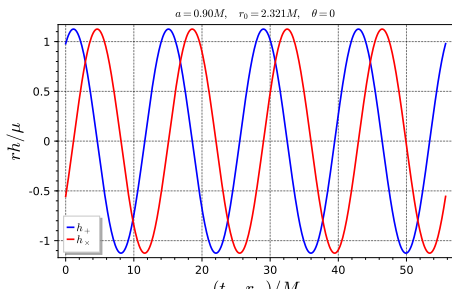
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Example for $a = 0.9 M$, $r_0 = r_{\text{ISCO}}(a)$ and viewing angle $\theta = 0$ (face-on)



Waveforms from circular orbits

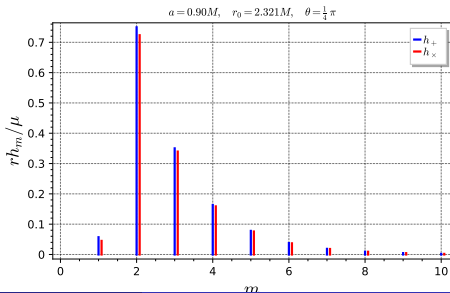
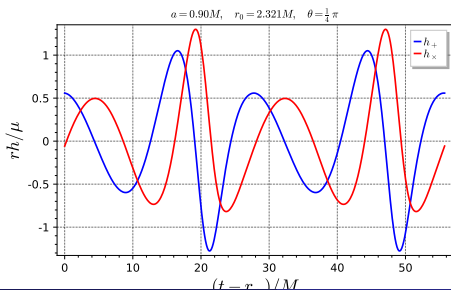
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Example for $a = 0.90M$, $r_0 = r_{\text{ISCO}}(a)$ and viewing angle $\theta = \pi/4$



Waveforms from circular orbits

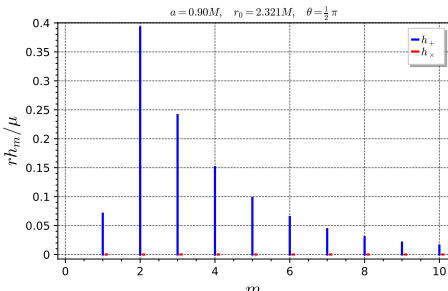
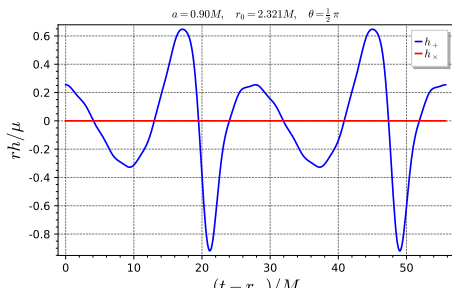
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Example for $a = 0.9M$, $r_0 = r_{\text{ISCO}}(a)$ and viewing angle $\theta = \pi/2$ (edge-on)



Implementation: the `kerrgeodesic_gw` package

All computations (GW waveforms, SNR in LISA, energy fluxes, inspiralling time, etc.) have been implemented as a **Python package** for the open-source mathematics software system **SageMath**:

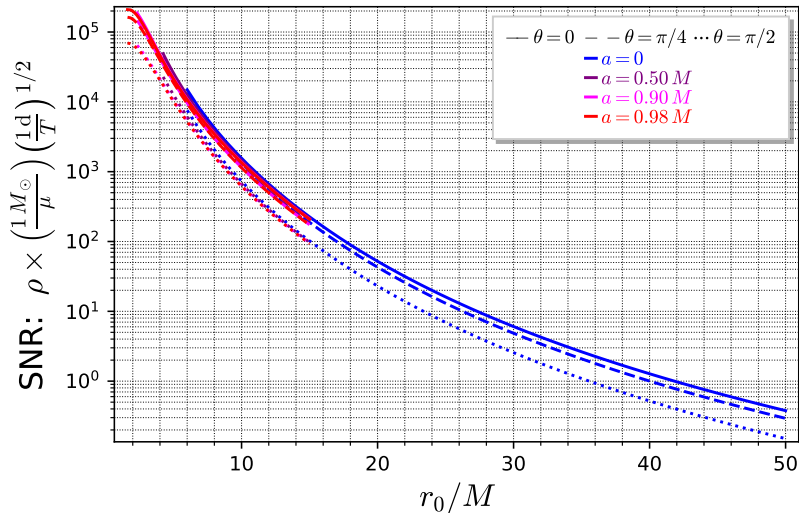
`kerrgeodesic_gw`

`kerrgeodesic_gw` is

- entirely **open-source**:
https://github.com/BlackHolePerturbationToolkit/kerrgeodesic_gw
- is distributed via **PyPi** (Python Package Index):
<https://pypi.org/project/kerrgeodesic-gw/>
- is part of the *Black Hole Perturbation Toolkit*:
<http://bhptoolkit.org/>

Signal-to-noise ratio in the LISA detector

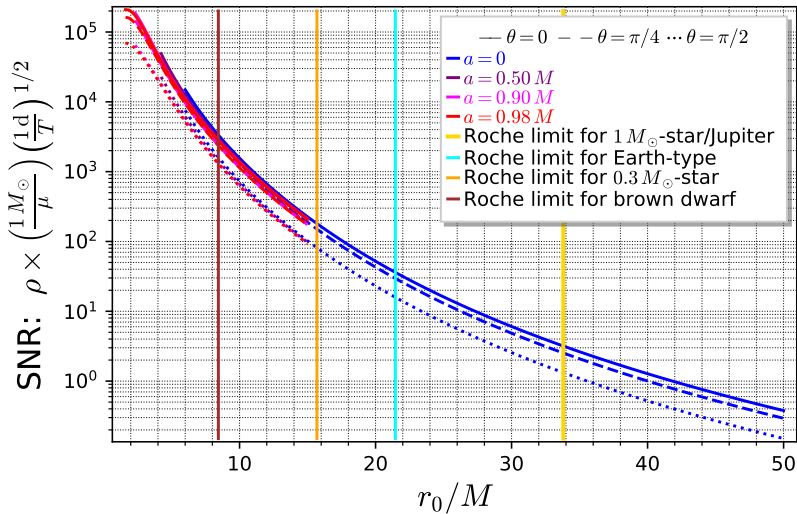
as a function of the circular orbit radius r_0



[Gourgoulhon, Le Tiec, Vincent & Warburton, A&A 627, A92 (2019)]

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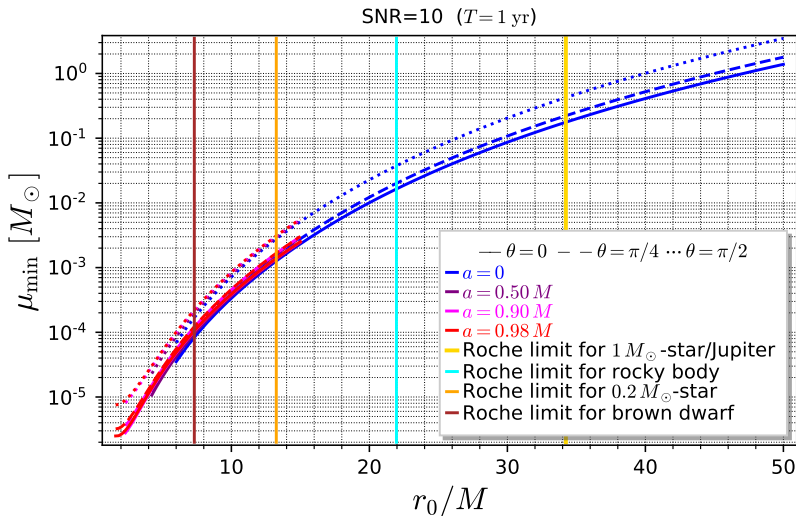


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Minimal detectable mass by LISA

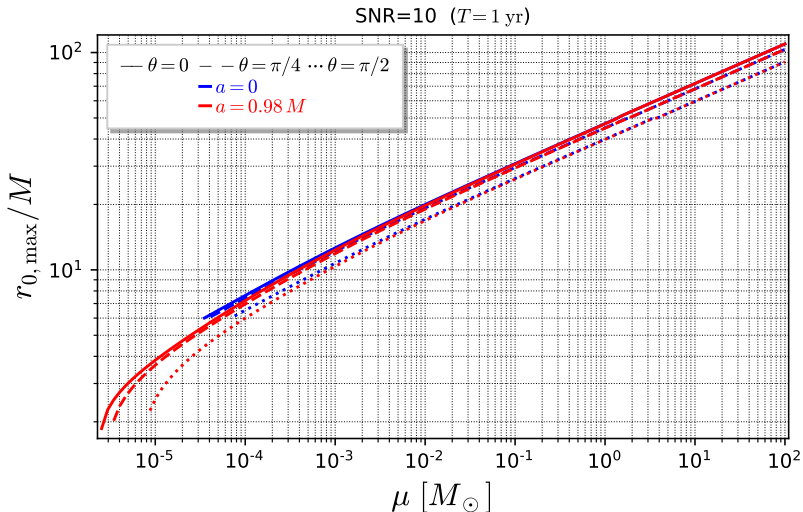
Detection criteria: $\text{SNR} \geq 10$

Observation time: $T = 1 \text{ yr}$



Maximum orbital radius for LISA detection

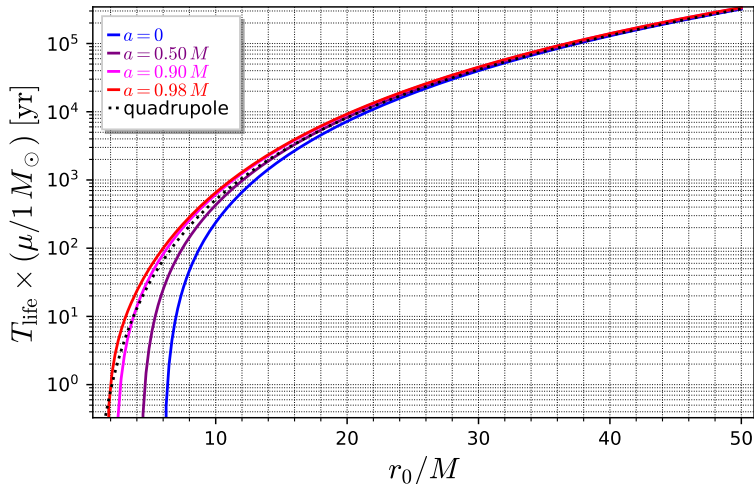
Maximum orbital radius $r_{0,\max}$ for a SNR = 10 detection by LISA in one year of data, as a function of the mass μ of the object orbiting around Sgr A*:



Detection probability governed by the life time of orbits

gravitational radiation reaction \implies slow inspiral motion

$T_{\text{life}}(r_0)$: time for a compact object to reach the ISCO starting from circular orbit of radius r_0



Time spent in LISA band

Inspiral time from orbit r_0 to orbit r_1 due to reaction to gravitational radiation:

$$T_{\text{ins}}(r_0, r_1) = \frac{M^2}{2\mu} \int_{r_1/M}^{r_0/M} \frac{1 - 6/x + 8\bar{a}/x^{3/2} - 3\bar{a}^2/x^2}{(1 - 3/x + 2\bar{a}/x^{3/2})^{3/2}} \frac{dx}{x^2(\tilde{L}_\infty(x) + \tilde{L}_H(x))}$$

where $\tilde{L}_{\infty, H}(x) := (M/\mu)^2 L_{\infty, H}(xM)$ and L_∞ (resp. L_H) is the total GW power emitted at infinity (resp. through the BH event horizon) by a particle of mass μ orbiting at $r = xM$

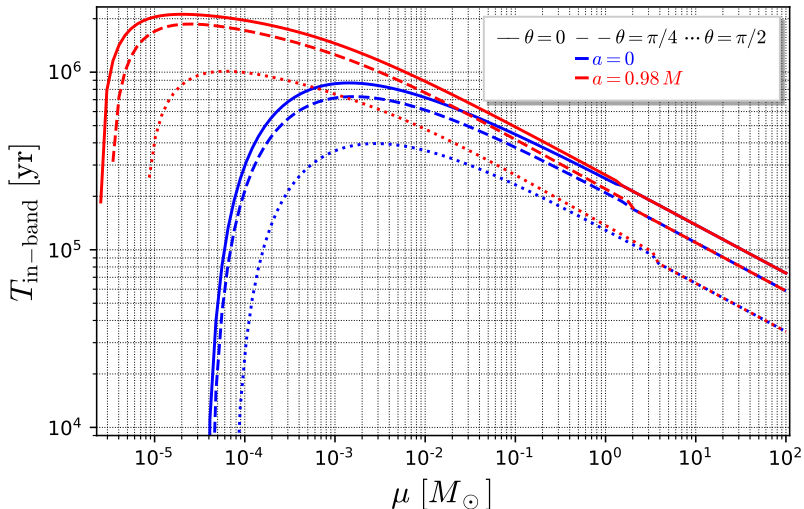
Compact object

$$T_{\text{in-band}} = T_{\text{ins}}(r_{0, \text{max}}, r_{\text{ISCO}}) = T_{\text{life}}(r_{0, \text{max}})$$

Main-sequence stars and brown dwarfs

$$T_{\text{in-band}} \geq T_{\text{in-band}}^{\text{ins}} = T_{\text{ins}}(r_{0, \text{max}}, r_{\text{Roche}})$$

Time in LISA band for an inspiralling compact object as a function of the compact object mass μ



Time in LISA band for brown dwarfs and main-sequence stars

Results for

- inclination angle $\theta = 0$
- BH spin $a = 0$ (outside parentheses) and $a = 0.98M$ (inside parentheses)

	brown dwarf	red dwarf	Sun-type	$2.4 M_{\odot}$ -star
μ/M_{\odot}	0.062	0.20	1	2.40
ρ/ρ_{\odot}	131.	18.8	1	0.367
$r_{0,\max}/M$	28.2 (28.0)	35.0 (34.9)	47.1 (47.0)	55.6 (55.6)
$f_{m=2}(r_{0,\max})$ [mHz]	0.105 (0.106)	0.076 (0.076)	0.049 (0.049)	0.038 (0.038)
r_{Roche}/M	7.31 (6.93)	13.3 (13.0)	34.2 (34.1)	47.6 (47.5)
$T_{\text{in-band}}^{\text{ins}} [10^5 \text{ yr}]$	4.98 (5.55)	3.72 (3.99)	1.83 (1.89)	0.938 (0.945)

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Brown dwarfs stay for $\sim 5 \times 10^5$ yr in LISA band

Conclusions

- GW emission and SNR in LISA for close circular orbits around Sgr A* has been computed in full general relativity.
- The time spent in LISA band ($\text{SNR} \geq 10$) during the slow inspiral has been evaluated.
- All computations have been implemented in the open-source SageMath package `kerrgeodesic_gw`, as part of the **Black Hole Perturbation Toolkit**.
- LISA has the capability to detect orbiting masses close to the ISCO as small as $\sim 10M_{\text{Earth}}$ or even $\sim 1M_{\text{Earth}}$ if Sgr A* is a fast rotator ($a \geq 0.9M$); this could involve primordial BHs or (hypothetical) very dense artificial objects.
- The longest times in-band, of the order of 10^6 years, are achieved for **primordial black holes** of mass $\sim 10^{-3}M_{\odot}$ down to $10^{-5}M_{\odot}$ (depending on Sgr A*'s spin), as well as for **brown dwarfs**, just followed by white dwarfs and low mass main-sequence stars.