

# PULSARS AND NEUTRON STARS: AN INTRODUCTION TO OBSERVATIONS AND CURRENT MODELS

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# Introduction

A star made of neutrons... is the main idea.

**BUT**

it is not the only feature:

- strong gravitational field
- electro-magnetic field
- transition to the surface and/or accretion of matter from the interstellar medium (see also lecture by N. Sandulescu)
- many possibilities for nuclear matter: protons, electrons, muons, pions,  $s$  quarks, hyperons, ...

⇒ self-gravitating object at nuclear density (not only neutrons)

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In 1932, British scientist James Chadwick discovers the existence of neutrons.

The legend says that, the very same evening, Lev Landau discusses about the possibility of dense stars made only of neutrons.



- density = nuclear density  $\rho_{\text{nuc}} \simeq 1.66 \times 10^{17} \text{ kg / m}^3$ ,
- let us take a mass  $M \simeq 1M_{\odot} = 2 \times 10^{30} \text{ kg}$ ,
- the radius should be  $R \simeq \sqrt[3]{\frac{3M}{4\pi\rho_{\text{nuc}}}} \simeq 14 \text{ km}$ ,
- and the surface gravity  $\sim 10^{11} \text{ m / s}^2$  (about  $10^9$  times that of the Sun).

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In 1967, at Cambridge University Antony Hewish lead an observational survey of extra-galactic radio-sources.

In august, his student Jocelyn Bell detects important signal fluctuations, which are observed to be periodic with a period of 1.337 s.



⇒ **pulsar** or *Pulsating Source of Radio* (PSR)

# PULSARS ARE NEUTRON STARS

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Originally, the sources were associated with pulsations of **neutron stars** or **white dwarfs** after thinking of little green men: the first pulsar has been called LGM before PSR B1919+21...

Thomas Gold, in 1968 identifies pulsars with rotating magnetized neutron stars and predicts a slight increase of their period due to energy loss.

⇒ with the discovery of the Crab pulsar (PSR B0531+21) with a period of 33 ms, white dwarfs are ruled out...





# HOW ARE THEY FORMED?

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In 1934, the astronomers Walter Baade and Fritz Zwicky make the hypothesis that *supernovæ* could represent the transition between main sequence (normal) stars and neutron stars.



- reasonable explanation for the source of energy of *supernovæ*
- in accordance with the discovery of Crab pulsar (*supernovæ* observed in 1054).

# STELLAR EVOLUTION

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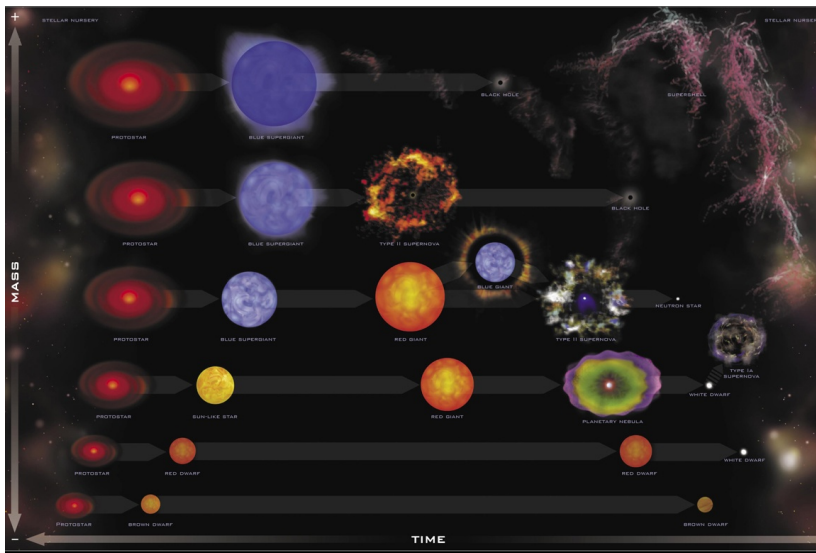
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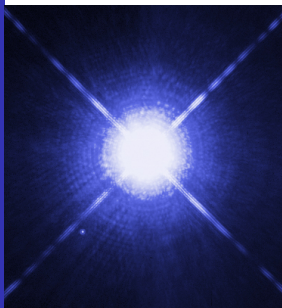
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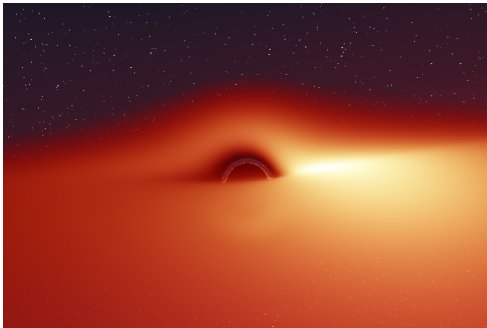




- white dwarfs have been observed as very hot (white) and very small (dwarf) stars;
- e.g. Sirius B has been observed with  $T_{\text{eff}} = 24000K$  and  $R \sim 5000 \text{ km}$ ;
- the density is therefore  $\rho \sim 10^9 \text{ kg/m}^3$ .

- gravity is balanced by degeneracy pressure of the electrons,
- if the mass is lower than some critical value, called *Chandrasekhar mass*  $\sim 1.5M_{\odot}$ .
- If they do not accrete matter, white dwarfs cool down and become fainter and fainter.

- a black hole is a region of space-time causally disconnected from asymptotic observers.
- the (geometric) boundary is called **event horizon**.



- inside the horizon, no light ray can reach any outside observer (black);
- what falls inside the horizon is definitely lost (hole) ...

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# Neutron star models

First one builds **equilibrium** models:

- assume that nuclear matter is a fluid;
- assume chemical equilibrium;
- give the gravitational law (self-gravitating body);
- write the hydrostatic equilibrium;
- give a law for the pressure as a function of nuclear matter density (temperature?)  $\Rightarrow$  **equation of state** (EOS).

The EOS specifies the nuclear matter properties and, in particular the strength of the strong interaction between particles, which is to equilibrate gravity.

Let us compare the escape velocity  $v_e$  from a self-gravitating body to the speed of light  $c$ :

$$\frac{1}{2}mv_e^2 = \frac{GMm}{R} \quad \text{and} \quad \frac{v_e^2}{c^2} = \frac{2GM}{Rc^2} = \Xi$$

$\Xi$  is called the **compactness** parameter of the body, and also measures the ratio between the gravitational potential energy and the mass energy. Its value is

- $10^{-6}$  for our Sun,
- $10^{-3}$  for a white dwarf,
- $0.4$  for a neutron star and
- $1$  for a black hole

⇒ it gives the influence of relativistic effects on the gravitational force: one needs **general relativity** (GR) to describe neutron stars!

# TOLMAN-OPPENHEIMER-VOLKOV SYSTEM

To describe the gravitational field in GR, one needs the **metric**  $g_{\alpha\beta}$  (static and spherically symmetric):

$$ds^2 = g_{\alpha\beta} dx^\alpha dx^\beta = -N^2 c^2 dt^2 + A^2 dr^2 + r^2 (d\theta^2 + \sin^2 \theta d\varphi^2)$$

The matter is assumed to be a perfect fluid and is given by its **stress-energy tensor**  $T^{\alpha\beta}$ :

$$T^{\alpha\beta} = \left( \rho + \frac{p}{c^2} \right) u^\alpha u^\beta + p g^{\alpha\beta}.$$

Defining  $m(r)$  and  $\Phi(r)$  from  $A = \left( 1 - \frac{2Gm}{rc^2} \right)^{-1}$  and

$N = \exp(\Phi/c^2)$ , Einstein and hydrostatic equations write:

$$\frac{dm}{dr} = 4\pi r^2 \rho$$

$$\frac{d\Phi}{dr} = \left( 1 - \frac{2Gm}{rc^2} \right)^{-1} \left( \frac{Gm}{r^2} + 4\pi G \frac{p}{c^2} r \right)$$

$$\frac{dp}{dr} = - \left( \rho + \frac{p}{c^2} \right) \frac{d\Phi}{dr}$$



In order to integrate the TOV system, one must specify an EOS:

- soon after their birth in *supernovæ*, neutron stars cool down below their Fermi temperature;
- temperature effects can be neglected and nuclear reactions are at equilibrium;
- **cold catalyzed matter at the endpoint of thermonuclear evolution.**

⇒ all state variables are functions of only one parameter, chosen to be e.g. the baryonic density  $n$ .

Defining the **adiabatic index**  $\Gamma(n) = \frac{n}{p} \frac{dp}{dn}$ , a family of simple EOSs, called polytropic EOSs, is obtained by assuming

$\Gamma = \text{const.}$

The TOV system is integrated specifying (in addition to the EOS):

- a value of the central density (to vary the resulting mass);
- regularity conditions at  $r = 0$ ;

Thus, the system can be integrated until  $p = 0$  at the surface ( $r = R$ ), where it is matched with vacuum spherical static solution (Birkhoff's theorem).

In GR, one must be cautious when defining global quantities:

- the **gravitational mass** can be defined for an isolated system, and here as  $M_g = \int_0^R 4\pi r^2 \rho(r) dr = m(R)$ .
- the **baryon mass** is given by the number of baryons contained in the star  $M_b = m_b \int_0^R 4\pi r^2 A(r) n(r) dr$ .
- the **gravitational redshift** is the frequency relative redshift undergone by a signal emitted at the surface of the star and measured by a distant observer

$$z = \left(1 - \frac{2GM}{Rc^2}\right)^{-1/2} - 1 = (1 - 2\Xi)^{-1/2} - 1.$$

# MAXIMAL MASS

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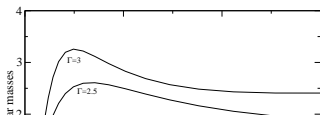
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Contrary to white dwarfs, where  $\Gamma \rightarrow 4/3$  near maximal mass, neutron star models exhibit a maximal mass as a general relativistic effect: more mass  $\Rightarrow$  more pressure to equilibrate  $\Rightarrow$  more gravity  $\Rightarrow$  more mass ...



Next step: take into account rotation; two possibilities

- Analytically perturb spherical models;
- Numerically compute full models in axisymmetry.

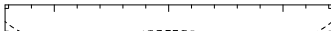
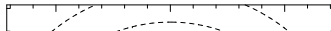
Assumption of **circularity** (no meridional convective currents),  
 $\Rightarrow$  four gravitational potentials, depending on  $(r, \theta)$ . Two more differences with spherical models:

- need to specify the **rotation law** as  $\Omega = f(r \sin \theta)$ ,  
 $f = \text{const}$  being rigid rotation;
- no more Birkhoff's theorem: the gravitational potentials must be integrated up to spatial infinity, where space-time is asymptotically flat.

With  $H = \log \left( \frac{e + p}{nm_b c^2} \right)$  the pseudo-enthalpy, the fluid equilibrium reads ( $\gamma$  is the fluid Lorentz factor):

$$H + \Phi - \log \gamma = \text{const.}$$

System of four coupled non-linear Poisson-like equations, with non-compact sources  $\Rightarrow$  use of computer techniques

Enthalpy isocontours ( $f=1$  kHz)Gravitational potential  $\Phi$  ( $f=1$  kHz)

# ROTATING MODELS: KEPLER LIMIT

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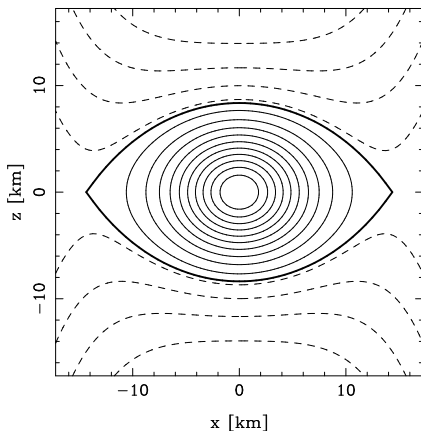
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In the case of rigid rotation, the angular frequency is physically limited by the **mass shedding limit**.

Enthalpy isocontours Keplerian frequency



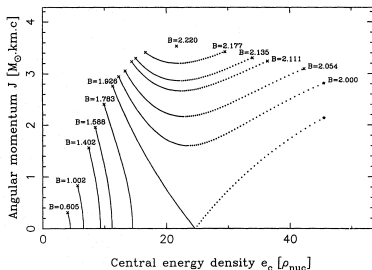
Also called **Kepler limit**,  
here  $\Omega_K \simeq 1100 Hz$ .

For each EOS one can  
define an absolute  
maximal rotational  
frequency.

For a given number of baryons, the effect of rotation is to:

- increase the radius of the star,
- decrease its central density,
- increase the gravitational mass.

The maximal mass associated to a given EOS is even more increased from the existence of **supermassive sequences**: rotating solutions that cannot exist in spherical symmetry.





Trying to improve the model...

- Neutrons are predicted to be superfluid;
- Other constituents are locked together by viscosity and magnetic field;  $\Rightarrow$  another fluid
- Both fluids are coupled by the strong nuclear force...  
 $\Rightarrow$  need for two-fluid framework, with non-dissipative coupling via entrainment and the nuclear “symmetry energy”.

# STRANGE QUARK STARS

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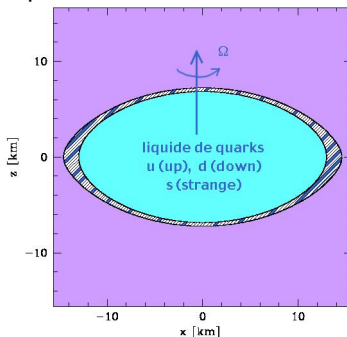
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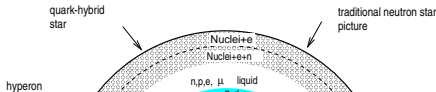
Deconfined beta-stable (u,d,s) quark matter has been proposed as the absolute ground state of matter at zero temperature.  
 ⇒some “neutron stars” could be “strange quark stars” ...  
 ... more compact, rotating faster, they have been invoked several times to explain anomalous observations.



⇒many numerical models to predict their exact properties.

Beside the “standard model” ( $n, p, e$  and  $\mu$ ) and the “strange quark star” picture, many possibilities have been considered for the inner core:

- pion condensate,
- hyperons,
- kaon condensate,
- quark condensate (color



# SOME POPULAR EQUATIONS OF STATE

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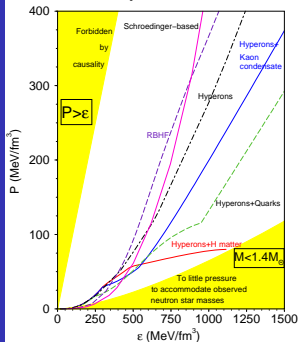
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EOS represents the major uncertainty when building models of neutron stars.  $\Rightarrow$  use of different EOSs present in the literature to compare all models to observations.



- polytropic EOSs,
- incompressible (constant density matter) or causal limit EOSs
- $(n, p, e$  and  $\mu)$  described by Skyrme potential (e.g. Skyrme Lyon 4 as in Douchin & Haensel 2001);
- MIT bag model for  $(u, d, s)$  quark matter.

$\Rightarrow$  very different conditions from laboratory experiments (high  $N/Z$ , low temperature, ...).

# STRUCTURE AND TRANSITION TO THE SURFACE

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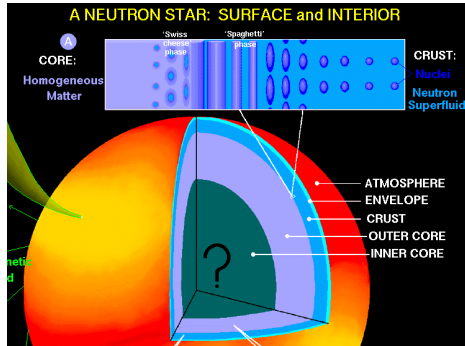
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The very interior is poorly known, but from the center, there must be a transition to neutron-rich fluid  $\Rightarrow$  neutrons dripping out of nuclei  $\Rightarrow$  crystal-like structure of nuclei for the outer crust.

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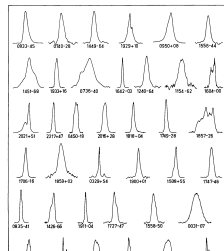
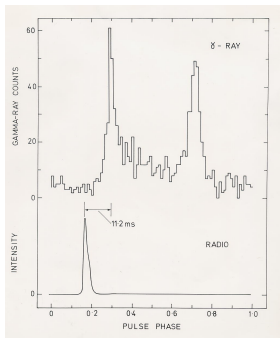
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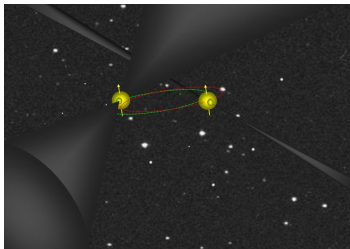
# Observations and constraints

Pulsars emit pulses in radio or X-rays, sometimes both (e.g. Vela):



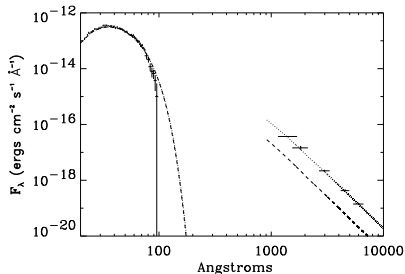
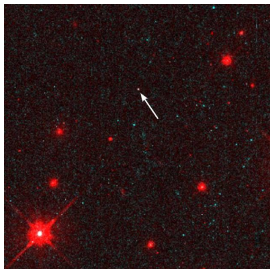
In 1974 R.A. Hulse and J.H. Taylor discovered first binary pulsar PSR B1913+16. Others have been discovered since then, including a double pulsar (both neutron stars are active and emit in our direction)

⇒ precise timing of binary pulsars have provided numerous valuable and very accurate observations



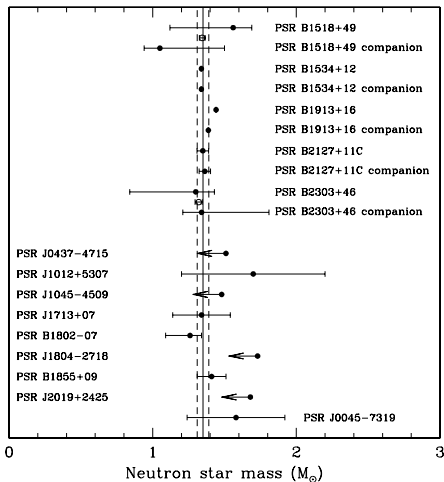


- in 1996, an isolated neutron star has been discovered in optics and X-rays;
- it shows no pulsar emission, only a surface thermal one, at  $T \sim 10^5 K$ ;
- called RX J185635-3754, it is situated at a distance of 120 pc from us (parallax measure by HST).



claimed to be a strange star for a while...

# MEASURE OF MASSES



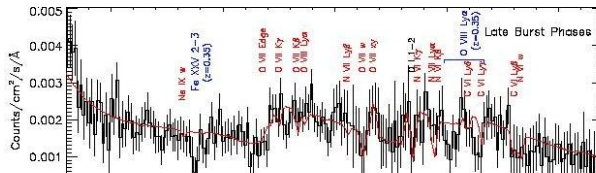
All measured neutron star masses have been obtained from binary systems:

- X-ray binaries (accreting neutron star);
- binary pulsars.

In the last case, the measures are very precise.

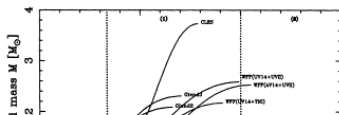
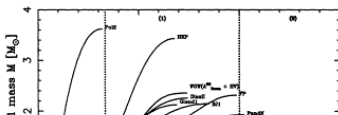
⇒ a realistic EOS must be able to give a maximal mass higher than  $\sim 1.6M_{\odot}$ .

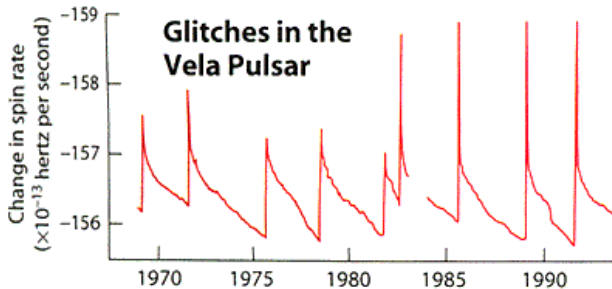
- Contrary to main sequence (normal) stars and white dwarfs, very few neutron stars radiate as black bodies...
- Strong magnetic field make atmosphere models complicated.  $\Rightarrow$  radii are not known!



# ROTATION PERIODS

- Rotation periods are easily deduced from pulse periods and range from  $P = 1.56 \text{ ms}$  to  $P = 8.5 \text{ s}$ .
- There has been a claim that a sub-millisecond ( $P = 0.5 \text{ ms}$ ) has been detected in the Large Magellanic Cloud, inside the SN1987a supernova, but withdrawn





Some pulsars exhibit sudden changes in the rotation period: instead of regularly slowing down, it shows rapid speed-up.

⇒ within the two-fluid framework:

- the outer crust (+fluid) is slowed down, not the inner fluid;
- until the stress (or interaction) between both becomes larger than some threshold.
- can turn into a “starquake”.

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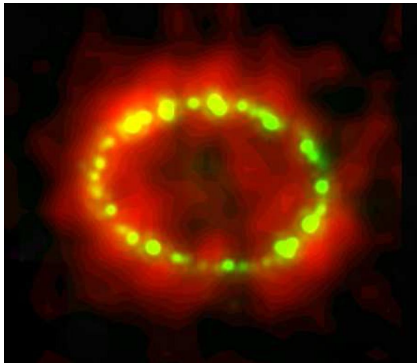
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One of the major astrophysical events of last century was the observation of the supernova SN1987a in the Large Magellanic Cloud and the detection of a dozen of neutrinos  $\nu_e$ .

These are an evidence of neutronization of matter through inverse  $\beta$ -reaction:



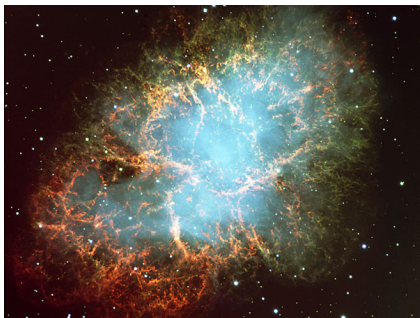
which is a confirmation of the scenario proposed by Baade and Zwicky 50 years earlier.

note: no pulsar has been ever found in the supernova remnant.

In 1969, the variation of the period of the Crab pulsar is measured giving

$$\dot{P} = 36 \text{ ns / day.}$$

Since  $\dot{P} > 0$ , the pulsar is slowing down, as predicted by Gold's model.



When estimating the age of a pulsar one can derive the simple formula

$$T \sim \frac{1}{2} \frac{P}{\dot{P}}.$$

If applied to the Crab (observed in 1054), one gets  $T \sim 1250$  years.

# NEUTRON STARS AS SOURCES OF GRAVITATIONAL WAVES

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Using a linearized version of Einstein equations describing the gravitational field from the matter-energy distribution:

The gravitational luminosity is

$$L \sim \frac{G}{c^5} s^2 \omega^6 M^2 R^4$$

( $s$  being a factor related to non-sphericity of the matter distribution); changing the formula to

$$L \sim \frac{c^5}{G} s^2 \left( \frac{2GM}{Rc^2} \right)^2 \left( \frac{v}{c} \right)^6$$

allows to see that good sources are

- non-spherical (and dynamically changing);
- compact ( $\Xi \sim 1$ );
- in relativistic motion.

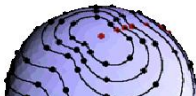
⇒ with the exception of first point, neutron stars are good potential sources...



To be efficient emitters, it is not enough to be flattened by rotation: they must have a variation in time of their **quadrupole** moment: e.g. deformation not symmetric / rotation axis.

Off-axis deformations can come from

- magnetic field;
- crust (1 mm high mountains);
- oscillations (vibrations) of the fluid.



Many modes, but most promising could be  $r$ -modes (Rossby waves) driven unstable by

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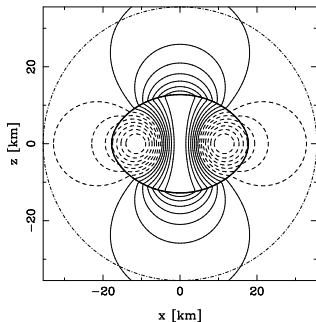
# Magnetic field and pulsar models

# MAGNETIZED NEUTRON STARS: NUMERICAL MODELS

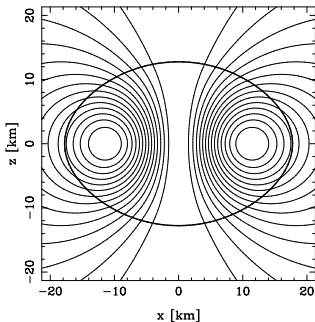
Magnetic field is important in the emission of pulsars

→ magnetized rotating models: aligned magnetic and rotation axis.

$A_t$  Potential



Magnetic field



In the system giving rotating models, one adds:

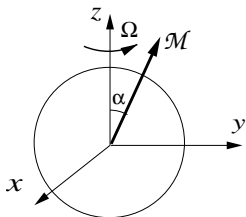
- two Poisson-like equations for the electro-magnetic potential (Maxwell equations)
- a magnetic potential in the equilibrium equation (“Bernoulli equation”).

To emit electro-magnetic waves, magnetic axis must not be aligned with the rotation one  $\Rightarrow$  oblique dipole (simplest analytic) model .

Accelerated dipole  $\Rightarrow$  radiation of its rotational energy  $E = \frac{1}{2} I \Omega^2$ :

$$\dot{E} = -\frac{2\pi}{3c^2\mu_0} \Omega^4 R^6 B_p^2 \sin^2 \alpha$$

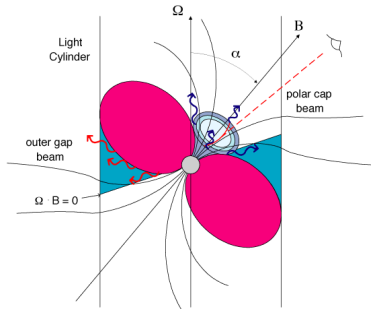
$\Rightarrow$  Deduce polar magnetic field  $B_p$  from observations of  $\dot{P}$ , assuming the star is a sphere.



For the Crab pulsar  $B_p \sin \alpha \simeq 5.3 \times 10^8$  T.

Open problem: the precise emission mechanism for pulsars.  
 Two main mechanisms are now proposed:

- The **polar cap** model: the emission region is near the magnetic pole, where the magnetic field is very strong, so pairs are created



- The **outer gap** model: the emission region is at the limit of closed magnetic field lines.

- If isolated rotating neutron stars are strong sources of gravitational radiation, then a part of pulsar slow-down should be due to this mechanism too. In general

$$\dot{\Omega} = -K\Omega^n,$$

where  $n$  is the **braking index** ( $n = 3$  for electro-magnetic waves and  $n = 5$  for gravitational waves).

- Observations (few and noisy) indicate  $n \lesssim 3$   
 $\Rightarrow$  electro-magnetic radiation seem to be at the origin of the pulsar slow-down.
- Usual pulsars have  $B_p \sim 10^{8-9}$  T,  $\Rightarrow$  depending on the exact current distribution inside the star, the magnetic deformation is rather small: in numerical models

$$\frac{p_{\text{mag}}}{p_{\text{fluid}}} \sim 10^{-9}.$$

# $(\dot{P}, P)$ DIAGRAM

Pulsars and neutron Stars

Jérôme Novak

Introduction

Models

Observations and constraints

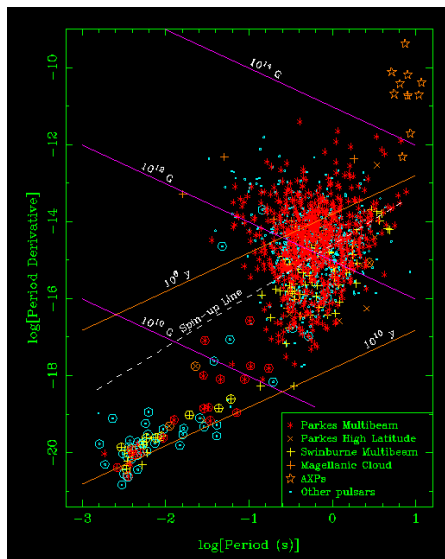
Pulsars

Laboratories for physics

The  $\sim 1500$  known pulsars can be displayed on a  $(P, \dot{P})$  plane, with the **death line** below which no radio emission can be seen.

Two categories of pulsars:

- “usual” ones, with  $P \sim 1$  s and  $B_p \sim 10^8$  T;
- **millisecond pulsars**, with much lower magnetic fields.



It is difficult to imagine the formation scenario of a 1 kHz rotating neutron star

⇒ millisecond pulsars must have been accelerated in some way.

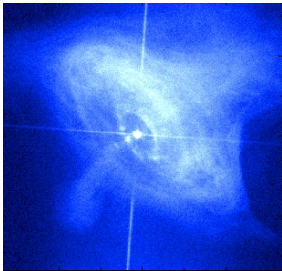
As nearly all millisecond pulsars are in binary systems, it seems plausible that the acceleration has occurred by accreting matter and angular momentum from the companion.



⇒ **pulsar recycling.**



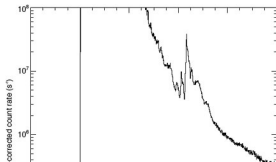
- Further out from the pulsar, a wind of relativistic particles is powered by the high-energy phenomena occurring in the pulsar magnetosphere.
- The wind is interacting with the supernova remnant and decelerated to sub-relativistic speed across a strong shock.
- These injected particles are thus radiating far from the pulsar in the **pulsar wind nebula** (also called “plerion”).
- The plerion can be observed at many wavelengths, but is mostly seen in X-rays. The typical plerion is the Crab



Some objects: Soft Gamma-ray Repeaters (SGR) and Anomalous X-ray Pulsars (AXP) observed in very-high energy emission, are associated with neutron stars with a very high magnetic field  $B_p \sim 10^{10-12}$  T.

⇒ magnetars

- These pulsars have high  $\dot{P}$  and low  $P$ ;
- they could represent as much as 10 % of all neutron stars (below death line);
- They can produce very strong bursts of X- and  $\gamma$ -rays, from the (glitch-like) re-arrangement of the crust, in which the magnetic field is pinned. The most striking case was the December 2004



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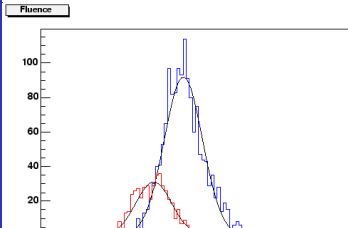
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# Neutron stars as probes for modern physics

From the 60s, military satellites have observed high-energy events producing bursts of  $\gamma$ -rays (GRB). Many dedicated satellites are now in operation, in conjunction with ground-based observations to catch the **afterglows**.

Two categories of GRBs have been observed:

- short ones, with a duration between 0.03 s and 2 s;
- long ones from 2 s to 1000 s.



# CENTRAL ENGINE FOR SHORT GAMMA-RAY BURSTS

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GRBs are the most energetic events, as the *supernovæ*. The emission mechanism is rather well understood, but the central powering engine is still uncertain.

For **short** GRBs, the standard model requires the merger of a binary system of neutron stars.

- The result of this collapse leads to the formation of a black hole surrounded by a hot torus emitting neutrinos.
- These could in turn deposit energy at the base of a relativistic jet,



# BINARY PULSARS IN GENERAL RELATIVITY

Pulsars and neutron Stars

Jérôme Novak

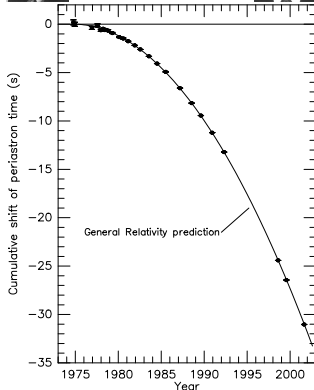
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Binary pulsars have been timed in an extreme accurate way, starting with PSR 1913+16, by Russel Hulse and Joseph Taylor.

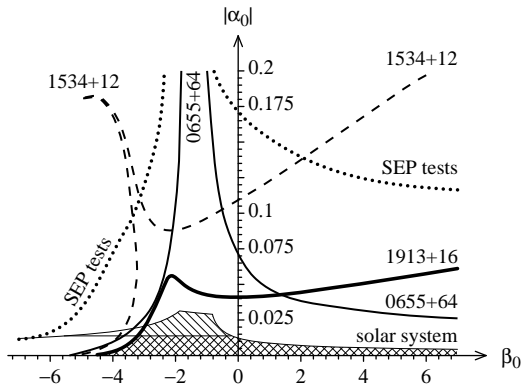
This precise timing allowed for a comparison with General Relativity prediction of

- the periastron shift,
- the angular momentum loss due to emission of gravitational waves.

Excellent agreement and check of General Relativity  $\Rightarrow$  Nobel Price in 1993

# TENSOR-SCALAR THEORIES OF GENERAL RELATIVITY

Other binary pulsars have been discovered and precisely timed. They have allowed for a comparison of general relativity with “more general” theories: **tensor-scalar theories of gravity**.



They are parameterized by two numbers  $(\alpha_0, \beta_0)$ ;  $\alpha_0 = 0$  and  $\beta_0 = 0$  correspond to general relativity.

# DETECTORS OF GRAVITATIONAL WAVES

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LIGO: USA, LOUISIANA



LIGO: USA, WASHINGTON



VIRGO: FRANCE/ITALY (PISA)



the arms of these  
Michelson-type LASERS  
are 3 km (VIRGO) and 4  
km (LIGO) long ... with  
almost perfect vacuum.  
⇒ they have just been  
built and are starting to  
acquire data.



# ASTROPHYSICAL SOURCES OF GRAVITATIONAL WAVES

## NEUTRON STARS AND BLACK HOLES

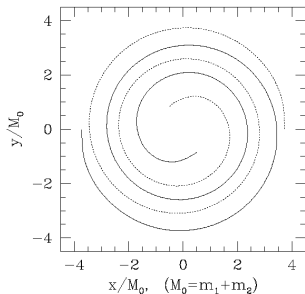
Many of astrophysical sources of high-frequency ( $10 \rightarrow 1000$  Hz) gravitational radiation imply neutron stars:

- at the very moment of their birth, *supernovæ* produce a fair amount of gravitational waves, although not very efficient (almost spherically symmetric);
- soon after that, *proto-neutron stars* can undergo growth of unstable oscillations, inducing deformations and a strong gravitational wave burst ;
- isolated *rotating neutron stars*: slightly deformed they emit gravitational waves. Not very efficient, but this mechanism is regular and the radiation can be integrated over a long period to increase the signal-to-noise ratio;
- *binary systems* of neutron stars or neutron star/black hole are losing angular momentum due to the emission of gravitational waves. They end up in a violent coalescence with a large amount of waves.

# ASTROPHYSICAL SOURCES OF GRAVITATIONAL WAVES

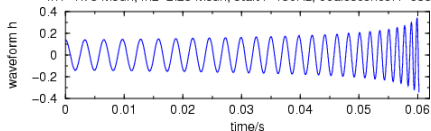
## WAVES

### BINARY NEUTRON STARS



### Gravitational Wave of Compact Binary Inspiral

$m_1=1.75$  Msun,  $m_2=2.25$  Msun, start  $f=150$ Hz, coalescence:  $f=635$ Hz



- Binary neutron stars and/or black holes are the best candidates to be detected by ground-based interferometric detectors of gravitational waves;
- fitting the quasi-periodic signal before the merger, one could recover the direction and the **masses**.
- the end of the signal is then sensitive to the **radius** of the neutron star and it might be possible (although difficult) to extract it/them from the signal.

Binary pulsars can give very precise measures of strong-field gravitational field.

⇒ isolated pulsars have been proposed as ultra-stable clocks.

If their physical properties are well understood, the long-term stability of pulsars as clocks is better than human-made atomic clocks.

One can see pulsars as precise clocks situated all over the universe and sending time signals to us...

⇒ In particular, **very low-frequency gravitational waves** ( $10^{-7} \rightarrow 10^{-4}$  Hz), it has been suggested that an extensive study and timing of pulsars could bring detections.

When building models of neutron stars and/or pulsar emission, the highest level of uncertainty comes from the EOS.

⇒ Could observations of neutron stars give any indication on the EOS of cold matter at nuclear density?

- a high mass or rotation rate can exclude some EOSs
- glitches can tell us about superfluidity
- a **real** thermal spectrum of a neutron star, together with its distance, could give the radius...
- high signal-to-noise ratio detection of gravitational waves from a binary neutron star could give masses and radii.

★ HOLY GRAIL ★

Get some points in the  
mass-radius diagram of  
neutron stars